

Kennicutt–Schmidt relation of galaxies over 13 billion years in the COLIBRE hydrodynamical simulations

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ABSTRACT

We investigate the correlation between star formation rate (SFR) surface density and gas surface density [known as the Kennicutt–Schmidt (KS) relation] at kiloparsec (kpc) scales across cosmic time ($0 \leq z \leq 8$) for galaxies with stellar masses $> 10^9 M_\odot$, using the COLIBRE state-of-the-art cosmological hydrodynamical simulations. These simulations feature on-the-fly non-equilibrium chemistry coupled to dust grain evolution and detailed radiative cooling down to ≈ 10 K, enabling direct predictions for the atomic (H I) and molecular (H₂) KS relations. At $z \approx 0$, COLIBRE reproduces the observed (spatially resolved) KS relations for H I and H₂, including the associated scatter, which we predict to be significantly correlated with stellar surface density, local specific SFR (sSFR), and gas metallicity. We show that the H I KS relation steepens for lower-mass galaxies, while the H₂ KS relation shifts to higher normalization in galaxies with higher sSFRs. The H₂ depletion time decreases by a factor of ≈ 20 from $z = 0$ to $z = 8$, primarily due to the decreasing gas-phase metallicity. This results in less H₂ and more H I being associated with a given SFR at higher redshift. We also find that galaxies with higher sSFRs have a larger molecular gas content and higher star formation efficiency per unit gas mass on kpc scales. The predicted evolution of the H₂ depletion time and its correlation with a galaxy’s sSFR agree remarkably well with observations in a wide redshift range, $0 \leq z \leq 5$.

Key words: galaxies: evolution – galaxy: formation – galaxies: ISM – galaxies: star formation.

1 INTRODUCTION

The connection between star formation and the interstellar medium (ISM) of galaxies is one of the most fundamental questions in extragalactic astronomy. Many surveys have been dedicated to understanding how the star formation rate (SFR) of galaxies, either in a global or in a spatially resolved manner, is

connected to the different phases of the ISM (see L. J. Tacconi, R. Genzel & A. Sternberg 2020; E. Schinnerer & A. K. Leroy 2024, for recent reviews).

In the local Universe the SFR correlates strongly with the molecular hydrogen (H₂) content, with the intrinsic scatter being significantly smaller than for the correlations obtained with either atomic hydrogen (H I) or total neutral hydrogen (H I + H₂) (e.g. A. K. Leroy et al. 2008; F. Bigiel et al. 2008; A. Schruba et al. 2011; I. Pessa et al. 2021). These relations have been studied at scales from a few hundred parsecs (pc) to a kilo pc (kpc) through

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correlating the surface densities of SFR, Σ_{SFR} , and gas, Σ_{gas} , in many tens of galaxies in the local Universe. This relation is known as the Kennicutt–Schmidt (KS) relation, named after the seminal works of M. Schmidt (1959) and R. C. Kennicutt (1983, 1989). Establishing how the KS relation evolves to higher redshift, z , for either H I or H₂, has been challenging, with the samples in which a resolved KS relation can be measured being limited to $\lesssim 10$ galaxies (e.g. J. Freundlich et al. 2013; C. E. Sharon et al. 2013, 2019; J. A. Hodge et al. 2015; D. Nagy et al. 2023; M. Béthermin et al. 2023; A. Zanella et al. 2024; J. S. Gómez et al. 2025). However, from global unresolved measurements of SFR and H₂ mass in thousands of galaxies, it has been possible to establish that the two quantities continue to correlate well up to at least $z \approx 5$ (e.g. L. J. Tacconi et al. 2020).

Local Universe observations have also helped establish that the KS relation, either in its atomic, molecular or total neutral gas form, displays a scatter that correlates with other galaxy properties; for instance, with the stellar surface density (e.g. Y. Shi et al. 2011, 2018; B. Dey et al. 2019; S. L. Ellison et al. 2020; L. Lin et al. 2022; J. Wang et al. 2024; T. Wong et al. 2024), suggesting that the stellar contribution to the local dynamical state plays a role in the gas-to-star conversion (e.g. L. Blitz & E. Rosolowsky 2006; E. C. Ostriker, C. F. McKee & A. K. Leroy 2010). Even when very dense gas traced by HCN (hydrogen cyanide) is used to measure the KS relation, observations continue to show scatter that correlates with other local properties (A. Usero et al. 2015; F. Bigiel et al. 2016), confirming that the underlying process of star formation is sensitive to local physical conditions, such as the stellar surface density (e.g. M. J. Gallagher et al. 2018; M. J. Jiménez-Donaire et al. 2019).

The physical drivers of the atomic and molecular KS relations are not understood. For instance, it is not clear to what extent molecular gas is a prerequisite for star formation as opposed to being a by-product of the cold and dense gas phase normally associated with star formation (J. Schaye 2004; S. C. O. Glover & P. C. Clark 2012). Several state-of-the-art cosmological galaxy formation models and simulations assume that H₂ is required to form stars (e.g. C. D. P. Lagos et al. 2011, 2018; R. S. Somerville, G. Popping & S. C. Trager 2015; C. G. Lacey et al. 2016; L. Xie et al. 2017; R. Davé et al. 2019; R. Feldmann et al. 2023; C. d. P. Lagos et al. 2024); while others only consider dense gas, whether it is molecular or not, and typically assume a fixed threshold density (and/or temperature) for star formation (e.g. M. Tremmel et al. 2017; A. Pillepich et al. 2018; V. A. Semenov, A. V. Kravtsov & N. Y. Gnedin 2018; O. Agertz et al. 2021). Going beyond a fixed density threshold for star formation, J. Schaye et al. (2015) attempted to encapsulate the transition from warm to cold ISM into a metallicity-dependent threshold density.

On smaller scales than usually probed by the KS relation results, observations of thousands of individual molecular clouds show very inefficient star formation, with only $\simeq 1$ per cent of the gas being converted into stars per dynamical time (see E. Schinnerer & A. K. Leroy 2024 for a review). This inefficiency is challenging to explain (C. Federrath 2015), with a combination of high turbulence (e.g. M. R. Krumholz & C. F. McKee 2005; C. Federrath & R. S. Klessen 2012) and efficient early stellar feedback (e.g. O. Agertz et al. 2013; Á. Segovia Otero et al. 2025) likely needed to cause a quick dispersal of clouds shortly after they have started to form stars.

It is still an open question how to combine the small scale, almost constant efficiency of gas-to-star conversion at the molecular cloud level, with the emergence of the atomic and molecular

KS power law on larger scales, at ≈ 1 kpc. Cosmological hydrodynamical simulations have recently attempted to shed light on this problem. K. Kraljic et al. (2024) used the NEWHORIZON cosmological hydrodynamical simulation box to study the emergence of the KS relation, focusing primarily on low-mass galaxies, which are more readily captured in their small cosmological volume of $(16 \text{ Mpc})^3$. They found the local gas turbulence to be a key driver of the scatter in the KS relation. This is perhaps unsurprising given the explicit dependence of the star formation efficiency on gas turbulence in their underlying star formation model. They also found that the observed slope of the KS relation at $z = 0$ is only established in their simulation at $z \approx 2\text{--}3$. More studies like this are needed to understand how general these results may be.

In this paper, we use a new suite of cosmological hydrodynamical simulations, COLIBRE (J. Schaye et al. 2026; E. Chaikin et al. 2026) to study the emergence of the atomic, molecular and total neutral gas KS relations, including their scatter. COLIBRE is especially well suited for this study as it includes a sophisticated model of the ISM that takes into account non-equilibrium chemistry of H and He coupled with a live dust model and gas cooling that tracks gas temperatures down to ≈ 10 K and accounts for self-shielding and local sources of radiation (S. Ploekinger et al. 2025; J. W. Trayford et al. 2026). This means that the H I and H₂ contents of the gas are *predictions* of the simulation. Previous generations of cosmological hydrodynamical simulations did not directly model the formation of H I and H₂ which could only be included in post-processing, imposing chemical and thermal equilibrium (e.g. C. d. P. Lagos et al. 2015; B. Diemer et al. 2017; R. Davé et al. 2020). A. J. Richings et al. (2022) show that these assumptions fail catastrophically at the high densities and low temperatures typical of the H₂-dominated gas, with differences up to an order of magnitude in the expected abundance of H I and H₂ in that regime. COLIBRE also includes a range of volumes and resolutions that allow us to probe a very large dynamic range in stellar mass, from dwarf galaxies to the most massive galaxies in the universe. COLIBRE uses a star formation model that assumes a fixed efficiency of gas conversion per free fall time of 1 per cent, and hence the emergence of molecular and atomic KS relations is not at all trivial.

This paper is organized as follows. Section 2 presents the COLIBRE simulations used in this work, briefly describing the physical models included. Section 2 also describes how we measure the KS relation and some of the other resolved properties of interest, such as gas metallicity and stellar surface density. Section 3 presents a detailed study of what COLIBRE predicts for the atomic, molecular and total neutral KS relations at $z = 0$, including correlations that arise between other resolved properties and the scatter of the KS relation. Section 3 also presents a comprehensive comparison with observations of the local Universe that probe the zero-point and scatter of the KS relation. Section 4 focuses on the predicted redshift evolution of the KS relation from $z = 8$ to $z = 0$, comparing with observations where possible, and exploring the correlation between the scatter of the KS relation and third parameters, such as gas metallicity and stellar surface density. Section 5 presents our conclusions. In Appendix A, we present extensive convergence tests of the KS relation.

2 THE COLIBRE SIMULATIONS

COLIBRE is a suite of cosmological hydrodynamical simulations of varying volume and resolution. A key aspect of COLIBRE is the leap in the sophistication of the modelling of several ‘subgrid’

modules compared to the previous generation of simulations in large cosmological volumes. Subgrid modules encapsulate the physics that happens below the resolution scale of the simulation and include (i) radiative cooling and heating; (ii) the formation, growth and destruction of dust grains; (iii) star formation from unstable gas; (iv) stellar mass loss and the rates of supernova (SN) Ia; (v) stellar winds, radiation pressure from starlight, and H II regions, i.e. stellar feedback processes that start earlier than SN feedback; (vi) SN feedback; (vii) black holes (BHs); (viii) active galactic nucleus (AGN) feedback; (ix) turbulent diffusion. Below we summarize the main characteristics of the subgrid physics modules. We refer to J. Schaye et al. (2026) for a full description:

(i) *Numerical setting.* The COLIBRE simulation suite uses the SWIFT code (M. Schaller et al. 2024) with the SPHENIX smoothed particle hydrodynamics (SPH) formulation for hydrodynamics (J. Borrow et al. 2022) and a fourth-order fast multipole method for gravity (W. Dehnen 2014). It includes both baryonic and dark matter (DM) particles, with four times more DM than baryonic particles, which helps suppress spurious transfer of energy from DM to stellar particles and hence improves the effective resolution for the baryonic component (A. D. Ludlow et al. 2019, 2021, 2023). Neutrino effects are modelled semi-linearly to account for their influence on structure growth (Y. Ali-Haïmoud & S. Bird 2013).

(ii) *Radiative cooling* is allowed down to ≈ 10 K. Cooling is computed using an updated version of the CHIMES chemical network (A. J. Richings, J. Schaye & B. D. Oppenheimer 2014a, b), integrated into the HYBRID-CHIMES framework (S. Ploekinger et al. 2025). The model treats non-equilibrium hydrogen and helium chemistry explicitly and uses equilibrium tables for metal cooling, corrected for non-equilibrium free-electron densities. It includes molecular hydrogen formation, photoionization, photodissociation, photoelectric heating by dust, and shielding by gas and dust, using a local Jeans column approximation (J. Schaye 2001; S. Ploekinger & J. Schaye 2020). Radiation sources include the cosmic microwave background, the UV/X-ray background (C.-A. Faucher-Giguère 2020), and a local interstellar radiation field (S. Ploekinger et al. 2025).

(iii) *Dust evolution* is modelled on-the-fly using six dust species (three compositions, two sizes), following J. W. Trayford et al. (2026). Dust grains are produced by AGB stars (F. Dell’Agli et al. 2017) and core-collapse supernovae (CCSNe; S. Zhukovska, H. P. Gail & M. Tieloff 2008), and grow via gas accretion following H. Hirashita & N. V. Voshchinnikov (2014). Destruction occurs via sputtering (J. C. Tsai & W. G. Mathews 1995), stellar feedback, and astration. Grain–grain collisions lead to shattering and coagulation (S. Aoyama et al. 2017; G. L. Granato et al. 2021). The dust physics is coupled to the chemistry and cooling modules.

(iv) *Star formation:* follows a gravitational instability criterion (F. S. J. Nobels et al. 2024), based on the local thermal and turbulent velocity dispersions. Specifically, gas particles are only eligible for star formation if they satisfy

$$\alpha \equiv \frac{\sigma_{\text{th}}^2 + \sigma_{\text{turb}}^2}{G \langle N_{\text{ngb}} \rangle^{2/3} m_{\text{g}}^{2/3} \rho_{\text{g}}^{2/3}} < 1, \quad (1)$$

where σ_{th} and σ_{turb} are the particle’s 3D thermal and turbulent velocity dispersions, respectively. N_{ngb} is the weighted mean number of neighbours in the SPH kernel, which is about 65 for a quartic spline SPH kernel adopted in COLIBRE. m_{g} is the mean, initial gas particle mass (which we list in Table 1 for the runs used

in this work). ρ_{g} is the gas density. σ_{th} and σ_{turb} are computed from a gas particle’s temperature and the local inter-particle velocity (peculiar plus Hubble) dispersion. Gas that meets the criterion in equation (1) is stochastically converted into stars using a M. Schmidt (1959) law ($\rho_{\text{SFR}} \propto \rho_{\text{gas}}/\tau_{\text{ff}}$) with a fixed efficiency of 1 per cent per free-fall time, τ_{ff} . The free-fall time-scales as $\rho_{\text{gas}}^{-0.5}$, so $\rho_{\text{SFR}} \propto \rho_{\text{gas}}^{3/2}$.

Because the instability criterion is evaluated at the resolution limit, the instability requirement in equation (1) enables higher-resolution runs to probe denser ISM conditions. This also means that as the resolution increases, a larger fraction of the SFR happens in denser gas compared to lower resolution runs (which we show in Section 3.1). F. S. J. Nobels et al. (2024), using a suite of idealized galaxy simulations, demonstrated that coarse-grained quantities, such as the observable KS law on kpc scales, can still converge with the resolution. We demonstrate that this is indeed the case using the fully cosmological simulations COLIBRE (see Appendix A1).

(v) *Chemical enrichment* models mass-loss from asymptotic giant branch (AGB) stars, core collapse SNe, and SNIa using updated nucleosynthetic yields (C. A. Correa et al. 2026). Tracked elements include the 11 cooling-dominant species (R. P. C. Wiersma et al. 2009), s-process elements (Ba, Sr), and r-process element (Eu). The SNIa rate follows an exponential delay-time distribution. The model uses metallicity-dependent stellar lifetimes (L. Portinari, C. Chiosi & A. Bressan 1998), with SNIa yields from the W7 model of S.-C. Leung & K. Nomoto (2019), core collapse SNe yields from K. Nomoto et al. (2006), K. Nomoto, C. Kobayashi & N. Tominaga (2013), and pre-SN winds from C. Kobayashi et al. (2006). C and Mg yields are scaled by 1.5 for massive stars to match APOGEE observations (H. Jönsson et al. 2018).

(vi) *Turbulent diffusion* is included to model small-scale mixing of metals and dust, following a diffusion approach motivated by Kolmogorov turbulence (F. J. Martínez-Serrano et al. 2008; T. H. Greif et al. 2009; S. Shen, J. Wadsley & G. Stinson 2010). The diffusion coefficient is tied to the local velocity shear, with the strength calibrated to reproduce Milky Way abundance patterns (C. A. Correa et al. 2026).

(vii) *Stellar feedback:* includes early feedback from stellar winds, radiation pressure, and H II regions (A. Benítez-Llambay et al. 2026). Core collapse SN feedback includes a stochastic thermal component (C. Dalla Vecchia & J. Schaye 2012, with updates) and a kinetic component to drive turbulence (E. Chaikin et al. 2023). SNIa feedback uses the same thermal model to prevent numerical overcooling.

(viii) *Black hole growth and AGN feedback:* The fiducial simulations use the black holes and thermal AGN feedback following C. M. Booth & J. Schaye (2009) and Y. M. Bahé et al. (2022), with modifications to improve the sampling of feedback from low-mass black holes. A subset of runs uses a hybrid AGN feedback model with both thermal and kinetic jets, and black hole spin tracking (F. Huško et al. 2026).

All feedback parameters are calibrated to reproduce the galaxy stellar mass function, galaxy sizes, and black hole masses at redshift zero, as described in E. Chaikin et al. (2026).

The adopted cosmology is consistent with the DES Y3 ‘3 × 2pt + All Ext.’ Λ CDM constraints (T. M. C. Abbott et al. 2022). This assumes a flat universe with key parameters: $H_0 = 68.1$ km/s/cMpc, $\Omega_{\text{m}} = 0.306$, $\Omega_{\text{b}} = 0.0486$, $\Omega_{\Lambda} = 0.693922$, $\sigma_8 = 0.807$, $n_s = 0.967$, $\Sigma m_{\nu} c^2 = 0.06$ eV

Table 1. COLIBRE simulations used in this paper. The columns list: (1) the name of the simulation, (2) comoving box size (L), (3) number of initial baryon (N_b) and (4) DM particles (N_{DM}), (5) initial mean particle masses of gas and (6) DM, (7) comoving gravitational softening length (ϵ_{com}) and (8) the maximum proper gravitational softening length (ϵ_{prop}). Here, ‘Thermal’ refers to runs using the thermal AGN feedback model, while ‘Hybrid’ refers to runs using the AGN feedback model of F. Huško et al. (2026). Our fiducial resolution is m6, and unless otherwise stated, L200m6 is the default simulation used for the analysis in this paper. To make this decision visually clear, we highlight the name of that simulation in bold.

(1) Name units	(2) L (cMpc)	(3) N_b	(4) N_{DM}	(5) Gas particle mass ($M_\odot$)	(6) DM particle mass ($M_\odot$)	(7) ϵ_{com} (ckpc)	(8) ϵ_{prop} (pkpc)
L025m5 (thermal)	25	752 ³	4×376^3	2.30×10^5	3.03×10^5	0.9	0.35
L025m6 (thermal)	25	376 ³	4×376^3	1.84×10^6	2.42×10^6	1.8	0.7
L025m7 (thermal)	25	188 ³	4×188^3	1.47×10^7	1.94×10^7	3.6	1.4
L050m6 (thermal)	50	752 ³	4×752^3	1.84×10^6	2.42×10^6	1.8	0.7
L050m6 (hybrid)	50	752 ³	4×752^3	1.84×10^6	2.42×10^6	1.8	0.7
L200m6 (thermal)	200	3008 ³	4×3008^3	1.84×10^6	2.42×10^6	1.8	0.7

(corresponding to one massive plus two massless neutrinos). The initial baryonic composition assumes $X = 0.756$ and $Y = 0.244$.

Structure identification proceeds in three stages:

(i) Friends-of-Friends (FoF): Identifies DM haloes using a standard linking length of 0.2 times the mean DM particle separation. Baryonic particles are linked to the nearest DM particle in a halo. Groups with < 32 particles are discarded.

(ii) HBT-HERONS (V. J. Forouhar Moreno et al. 2025): Tracks subhaloes across snapshots using a hierarchical bound-tracing approach, improving on HBT+ (J. Han et al. 2012, 2018). It excels at identifying satellite subhaloes near halo centres and avoids issues like artificial mass fluctuations at pericentres (V. J. Forouhar Moreno et al. 2025), which can affect the reconstruction of merger trees (Á. Chandro-Gómez et al. 2025).

(iii) SOAP (R. McGibbon et al. 2025): Calculates halo and galaxy properties within various 3D and projected apertures. The default aperture for galaxy properties used in this work is a sphere of radius 50 proper kpc (pkpc), centred on the most bound particle. Spherical overdensity masses are also computed.

The combination of FoF, HBT-HERONS, and SOAP provides a detailed and robust catalogue of halo and galaxy properties, that enables a consistent post-processing of COLIBRE simulation data. From the SOAP catalogues, we use stellar masses, instantaneous SFRs, H I, and H₂ masses. We also use these properties to define the specific SFR, $sSFR \equiv SFR/M_*$.

In this study, we select all galaxies with $M_* > 10^9 M_\odot$ and $SFR > 0$ to construct their resolved KS relation. We apply the same selection to snapshots from $z = 0$ to $z = 8$. This selection yields 106 572 galaxies at $z = 0$ in the L200m6 box, with the number of galaxies increasing to 120 398 at $z \approx 0.5$, followed by a steady decrease to 289 at $z = 8$. Appendix A1 presents a convergence study that explains our reasoning for choosing $M_* > 10^9 M_\odot$ as our minimum stellar mass for this analysis.

We make reference throughout this work to gas particles belonging to the hot ionized medium (HIM), warm neutral medium (WNM), and cold neutral medium (CNM). For simplicity, these phases are defined purely by gas temperature, acknowledging that this classification does not uniquely map on to ionization state (e.g. the WNM may contain ionized gas and the HIM may contain neutral gas; see Fig. 3). Specifically, we define the phases as follows:

$$\begin{aligned}
 \text{HIM} : T &\geq 10^{4.5} \text{ K}, \\
 \text{WNM} : 10^3 \text{ K} &< T < 10^{4.5} \text{ K}, \\
 \text{CNM} : T &\leq 10^3 \text{ K}.
 \end{aligned} \tag{2}$$

2.1 COLIBRE runs used in this work

Table 1 presents the COLIBRE runs used in this work. Throughout this paper, we will be primarily using the L200m6 run (using a cubic volume of 200 comoving Mpc on a side, and an initial mean particle mass of $10^6 M_\odot$). When extending the lower stellar mass limit from $10^9 M_\odot$ to $10^8 M_\odot$, we also employ the higher resolution run L025m5 (a cubic volume of 25 comoving Mpc on a side, and an initial mean particle mass of $10^5 M_\odot$). In Appendix A1, we use the L025 runs of different resolutions to study how well converged the predicted KS relations are with the numerical resolution. Throughout this work, we primarily use the ‘Thermal AGN feedback’ runs of COLIBRE. J. Schaye et al. (2026) present the full suite of COLIBRE runs, which include more boxes and variations with the ‘hybrid AGN feedback’ model. AGN feedback in COLIBRE primarily affects massive galaxies, and the varying AGN feedback models have the greatest impact on the population of massive, quenching/quenched galaxies. We confirm in Appendix A2 that the Thermal and Hybrid AGN feedback models predict KS relations at $z = 0$ that are almost indistinguishable from each other. Thus, throughout this paper we only analyse runs that adopt the Thermal AGN feedback model.

2.2 Constructing resolved maps

We implement several different ways of computing resolved maps of H I, H₂, SFR, stellar mass, gas metallicity, and gas velocity dispersion, which we describe and compare in Appendix A3. This is done for the purpose of testing the dependence of the KS relation on the way in which radial profiles are computed. We find very good agreement of the KS relation among different methods, and hence we limit ourselves to using the one described below.

Annuli on face-on galaxy: We first orient galaxies to be face-on using their stellar angular momentum vector. The latter is computed using all bound stellar particles within the 50 pkpc spherical aperture, centred on the most bound particle. We then take concentric rings of width Δr , and those are extracted out to 50 pkpc. From the particles belonging to each ring, we measure the H I, H₂, stellar mass, SFR and dust content of the ring, by summing their masses and SFRs. For the total dust mass, we sum the masses in the 6 dust species. For the metallicity, we compute the total number of oxygen atoms and divide by the total number of hydrogen atoms associated with gas particles in the ring that have a temperature $< 10^{4.5} \text{ K}$ and $> 0.1 \text{ cm}^{-3}$, respectively, excluding metals locked in dust. We then multiply by a factor of 12 to obtain

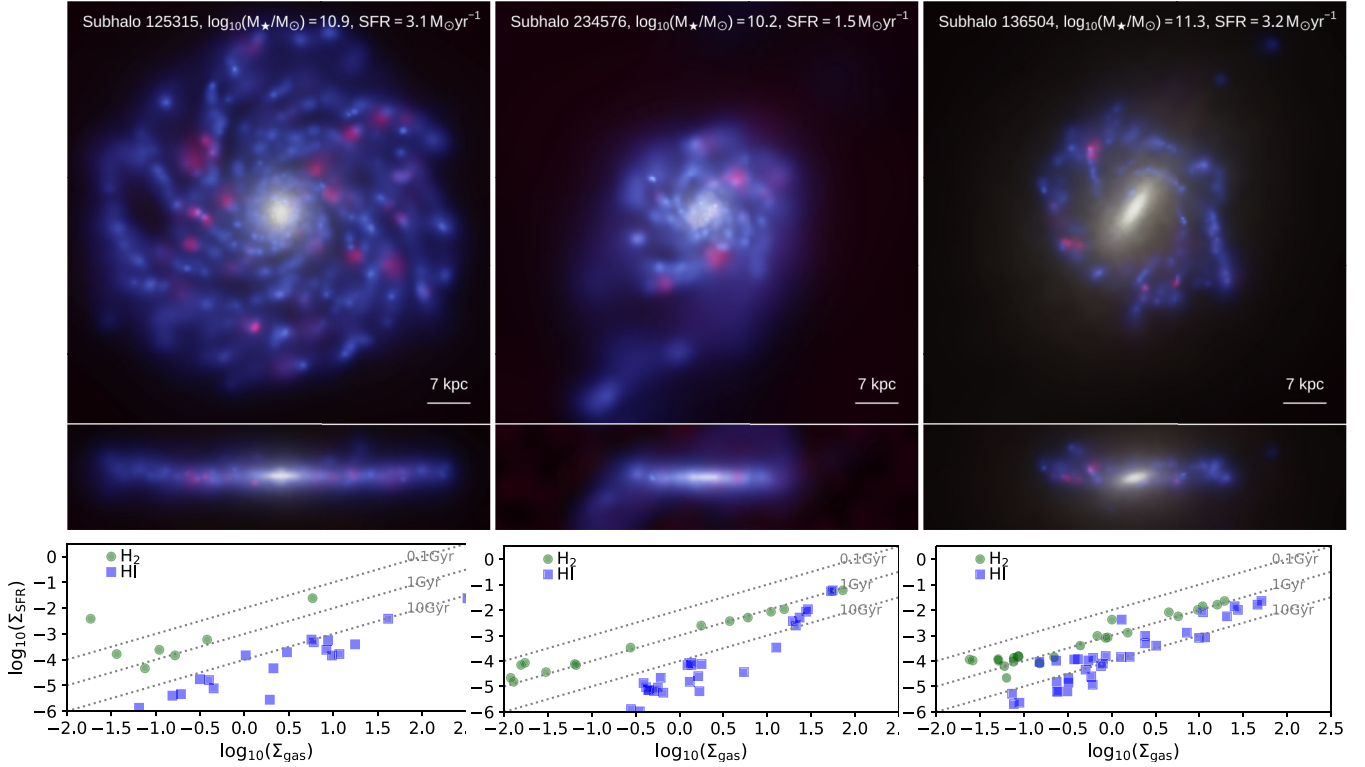


Figure 1. Three example galaxies at $z = 0$ in the L025m6 COLIBRE simulation used in this work. The top and middle panels show face-on and edge-on views of the gas and stellar disc, which were produced using R-package SWIFT. The gas is coloured by the ionization fraction, with blue indicating neutral gas ($\text{H I} + \text{H}_2$), and red ionized gas (primarily H II regions when the red compact regions are in the disc). The stellar component is shown with a gold colour, with the intensity of the colour scaling on the stellar mass density. Each figure shows at the top the stellar mass and SFR of the subhalo and the subhalo index within the SOAP catalogue. We also indicate with a white bar a physical scale for reference. The bottom panels show the individual H I (blue squares) and H_2 (green circles) KS tracks of the galaxies at the top. The tracks are computed using our fiducial ANNULI-FACE method. In the bottom panels, Σ_{SFR} and Σ_{gas} are in units of $\text{M}_\odot \text{ yr}^{-1} \text{ kpc}^{-2}$ and $\text{M}_\odot \text{ pc}^{-2}$, respectively. The dotted lines indicate fixed depletion times of 0.1, 1, and 10 Gyr as labelled. The left and middle galaxies are examples of normal star-forming galaxies, while the one on the right is an example of a strong bar galaxy that has emptied most of its inner region.

the metallicity in observer units of $12 + \log_{10}(\text{O}/\text{H})$. We also save the number of particles that are used for these calculations in each ring and include only annuli with a number of gas particles larger than N_{min} .

Our fiducial parameters are $\Delta r = 1$ proper kpc and $N_{\text{min}} = 10$ particles. Appendix A3 shows that our results are insensitive to the precise values if they are within $\Delta r = 0.5\text{--}2$ kpc and $N_{\text{min}} = 1\text{--}100$. Note that here we do not place constraints on the vertical distance from the disc plane and simply take all the particles that are bound to the subhalo. To turn masses into surface densities, we simply divide them by the area of the ring. We refer to this method as ANNULI-FACE in Appendix A3 when comparing with other methods.

The resolved maps above are computed for every galaxy in the L200m6 volume with $M_* > 10^9 \text{ M}_\odot$ and $\text{SFR} > 0$ at different snapshots, but because we place a N_{min} threshold for the minimum number of gas particles per ring, some rings in galaxies do not contribute to the measurement of the KS relation as they are too gas poor.

With these resolved maps, we define the surface densities of SFR (Σ_{SFR}), stellar mass (Σ_*), H I (Σ_{HI}), H_2 (Σ_{H_2}), total neutral gas ($\Sigma_{\text{HI}+\text{H}_2}$), dust (Σ_{dust}), the local gas metallicity $12 + \log_{10}(\text{O}/\text{H})$ and local specific SFR ($s\Sigma_{\text{SFR}} = \Sigma_{\text{SFR}}/\Sigma_*$). We also compute the cool gas vertical velocity dispersion profile, by selecting only the gas particles with a temperature $T < 10^{4.5} \text{ K}$ in each annulus and

calculating their mass-weighted velocity dispersion:

$$\sigma_{\text{cool}} = \sqrt{\frac{\sum (v_{z,i}^2 + \sigma_{\text{th}}^2) m_i}{\sum m_i}}, \quad (3)$$

where $v_{z,i}$ is the velocity component of gas particle i parallel to the stellar angular momentum vector in the centre-of-mass rest frame of the galaxy, $\sigma_{\text{th}} = \sqrt{k_B T / \mu m_H}$ is the one-dimensional thermal velocity dispersion, m_i is its mass, T its temperature, k_B is Boltzmann's constant, μ is the atomic weight and m_H is the hydrogen atom's mass. We tested including limits on the distance of the gas particles to the mid-plane of the galaxy disc, following the method presented in E. Jiménez et al. (2023) for the EAGLE simulations, but we do not see a noticeable effect on our results. For simplicity, we thus do not include a limit on distance to the mid-plane.

Fig. 1 shows examples of $z = 0$ galaxies in the L025m6 simulation with their corresponding H I and H_2 KS relations. Face-on and edge-on images (top and middle panels) were built with R-package SWIFT.¹ Individual symbols in the bottom panels correspond to individual rings in the example galaxies using the ANNULI-FACE method to measure the KS relation. In Fig. B1, we show additional examples of star-forming and passive galaxies,

¹The code is publicly available at <https://github.com/obreschkow/swift>.

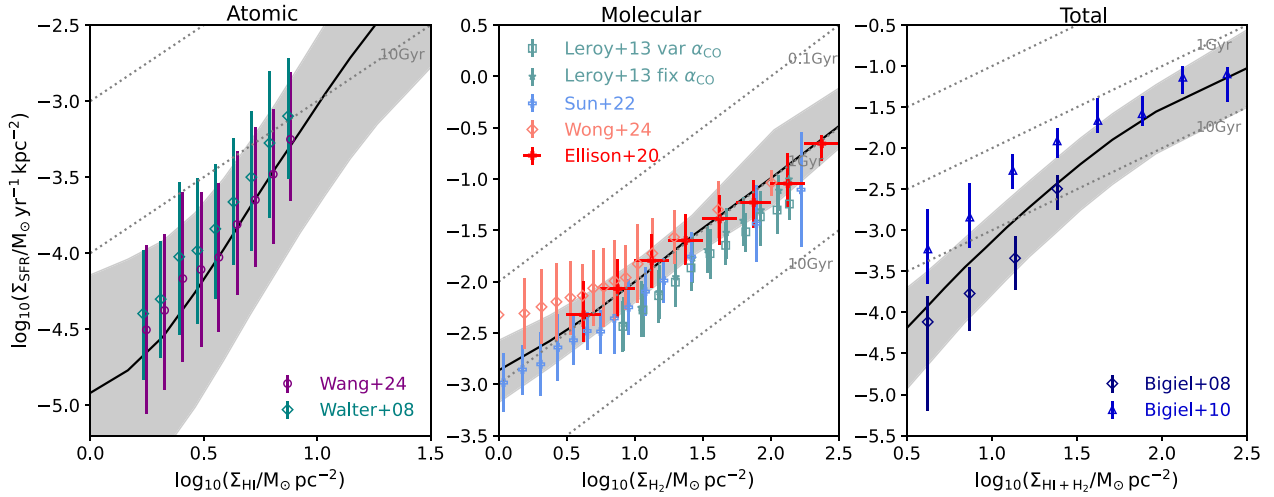


Figure 2. KS relation for H I (left panels), H₂ (middle panels) and total neutral hydrogen (H I + H₂) (right panels) for galaxies at $z = 0$ with $M_{\star} \geq 10^9 M_{\odot}$ in the COLIBRE simulation L200m6 (see Table 1). The KS relations are calculated as azimuthal averages in rings with a width of 1 kpc. Black lines and shaded regions show the median and 16th–84th percentile ranges, respectively, for COLIBRE. The bins are chosen to have a similar number of rings per bin, and always ≥ 10 . For reference we show in each panel dotted lines of constant gas depletion time (0.1, 1, and 10 Gyr), as labelled. We show observations from F. Walter et al. (2008), J. Wang et al. (2024) for H I, A. K. Leroy et al. (2013), S. L. Ellison et al. (2020), J. Sun et al. (2022), T. Wong et al. (2024) for H₂ and F. Bigiel et al. (2008), F. Bigiel et al. (2010) for total neutral gas, as labelled in each panel. For A. K. Leroy et al. (2013) we show two sets of values, using a constant and a variable CO-to-H₂ conversion factor (α_{CO}) to demonstrate the typical level of systematic uncertainty introduced by this conversion. Note that the x - and y -axes ranges are different in each panel and chosen to span the range over which observations are reported.

with the latter still having some low-level SFR. In general, we find that in a single galaxy, the H I KS relation has a lower amplitude and a similar or steeper slope than the H₂ KS relation. We will quantify these trends in Section 3.

We measure the correlation between Σ_{SFR} and other spatially resolved properties with the deviations from the primary KS relations. The idea is to understand how the scatter of the KS relation correlates with other resolved properties and to pinpoint the strongest correlations. This has been done in the past in observations to tease out the dependence of the scatter of one resolved scaling relation on another resolved property (e.g. S. L. Ellison et al. 2020). With this in mind, we define:

(i) Δ SFE: we fit power laws to the relations between Σ_{SFR} and Σ_{HI} , Σ_{H_2} or $\Sigma_{\text{HI}+\text{H}_2}$, using a χ^2 -minimization in log-space. We refer to these fits as KS_{HI} , KS_{H_2} and $\text{KS}_{\text{HI}+\text{H}_2}$. Using these fits, we then define $\Delta \text{SFE} = \log_{10}(\Sigma_{\text{SFR}}) - \text{KS}_x$, with x being H I, H₂, or H I + H₂ (i.e. deviations from the primary relations).

(ii) Δ SFR – prop: characterizes the correlations between the deviations from the main relation between $\log_{10}(\Sigma_{\text{SFR}})$ and another property. For our purpose, the properties we study are: $\log_{10}(\Sigma_{\star})$, $\log_{10}(\Sigma_{\text{dust}})$, and $\log_{10}(\text{O}/\text{H})$. We also use a linear fit to characterize all these correlations.

Appendix A1 presents a convergence study of the KS relation for H I, H₂, and H I + H₂, using the L025m5, L025m6, L025m7 simulations. For the H₂ KS relation, we see convergence at higher gas surface densities only, $\gtrsim 3 M_{\odot} \text{pc}^{-2}$. Comparing L025m5 and L025m6, we see that the H₂ KS relation converges at $\Sigma_{\text{H}_2} \gtrsim 1 M_{\odot} \text{pc}^{-2}$ for the m6 resolution at $z = 0$. The surface density at which the H₂ KS relation at m6 resolution appears to be converged decreases to $0.1 M_{\odot} \text{pc}^{-2}$ at $z = 2$. The median H I KS relation is very well converged at $z = 0$ at all gas surface densities, whereas the scatter tends to increase with increasing resolution. At $z = 2$, however, the H I KS relation at m6 starts to deviate from the results at m5 at $\lesssim 10 M_{\odot} \text{pc}^{-2}$. We will use these surface

density thresholds when fitting the KS relation of galaxies in the L200m6 box. We also show that selecting a minimum number of particles per annulus of 10, 50, or 100 has little impact on the resulting KS relation. Unless otherwise stated, throughout this paper, we require a minimum number of gas and stellar particles per annuli of 10.

3 KENNICUTT-SCHMIDT RELATION AT $z = 0$

In this section, we explore what COLIBRE predicts for the KS relation at $z = 0$ and present a comprehensive comparison with observations. Where necessary, we subtract the contribution of helium from the observational data (typically applied as a constant multiplicative factor) to compare with our predicted surface densities of H I, H₂, and H I + H₂. We also convert from a E. E. Salpeter (1955) or P. Kroupa (2001) IMF to our adopted G. Chabrier (2003) IMF where necessary. We do this following the conversion factors between different SFR tracers from different IMFs tabulated in Appendix A of V. Gonzalez-Perez et al. (2014). For stellar masses, we divide the masses derived using an E. E. Salpeter (1955) and P. Kroupa (2001) IMF by 2 and 1.14, respectively, to convert them to a G. Chabrier (2003) IMF.

3.1 KS relation: comparison with observations and characterization

We first aim to establish how well COLIBRE matches the local Universe observations of the KS relation. Fig. 2 shows the KS relation at $z = 0$ predicted by COLIBRE, separating the dependence on H I, H₂ and total neutral hydrogen. Here we use the *annuli on face-on galaxy* method. Each galaxy contributes up to 50 data points (i.e. rings) to each plane. However, because we set a minimum number of gas particles per ring to consider it in our analysis, we find that at $z = 0$ the median number of rings contributed to the measurement of the KS relation per galaxy is

only 7. Observational data from F. Walter et al. (2008), F. Bigiel et al. (2008), F. Bigiel et al. (2010), J. Wang et al. (2024), and S. L. Ellison et al. (2020) correspond to pixel-wise measurements of sizes 1, 0.75, 0.6, 1, and 0.3 kpc, respectively. Data from A. K. Leroy et al. (2013), J. Sun et al. (2022), and T. Wong et al. (2024) correspond to annuli of width 1 kpc, which exactly mimics our method. We note, however, that F. Bigiel et al. (2008) demonstrated that for individual galaxies measurements of the KS relation in pixels versus radial annuli returned indistinguishable results. This agrees with what we show for COLIBRE in Appendix A3.

COLIBRE is in remarkable agreement with observations for all three gas phases. For H I, COLIBRE tends to prefer a lower amplitude, very close to what J. Wang et al. (2024) reported. The difference between F. Walter et al. (2008) and J. Wang et al. (2024) is that the latter includes the more diffuse H I which tends to be missed by interferometric surveys. For H₂, we find that the COLIBRE amplitude is closest to the observational constraints from S. L. Ellison et al. (2020) and T. Wong et al. (2024), which correspond to a depletion time of or slightly above 1 Gyr. We note however, that the differences between different observations are within the typical systematic uncertainties associated with the CO-to-H₂ conversion factor (α_{CO}). For the total neutral hydrogen, we find that COLIBRE's predictions are in between the observational estimates at surface densities $\lesssim 10^{1.5} \text{ M}_{\odot} \text{ pc}^{-2}$, while at higher densities, there is agreement with F. Bigiel et al. (2010), which is the only data set available at those surface densities.

The agreement with observations was not guaranteed. In COLIBRE, the only free parameter in the star formation model is the fraction of gas mass converted into stars per free-fall time, provided the gas satisfies the instability criterion of equation (1). The adopted value is taken from observations of local universe molecular clouds (e.g. E. Schinnerer & A. K. Leroy 2024) rather than to ensure agreement with the observed KS relation, and it does not know how much H I or H₂ a gas particle has. The latter can then lead to large variations in the KS law with gas metallicity for example, which we return to in Section 3.2.1.

The predictions of COLIBRE presented here agree with those in F. S. J. Nobels et al. (2024) at $\log_{10}(\Sigma_{\text{H}_2}/\text{M}_{\odot} \text{ pc}^{-2}) \gtrsim 0.5$. At lower values, the simulation results in F. S. J. Nobels et al. (2024) drop towards longer depletion times. COLIBRE instead tends to continue on the constant depletion time line or even flatten a bit towards shorter depletion times. This difference is not necessarily surprising as there are many components to the physical model of COLIBRE that were not present in F. S. J. Nobels et al. (2024), including the non-equilibrium chemistry component, live dust and chemical evolution model and the overall cosmological context, which is lacking in F. S. J. Nobels et al. (2024). Regarding the non-equilibrium chemistry, by applying CHIMES to the FIRE simulations, A. J. Richings et al. (2022) show that explicitly solving for the chemistry of the gas leads to varying H I and H₂ mass fractions at low densities ($n_{\text{H}} \lesssim 1 \text{ cm}^{-3}$), which will then change the relation between SFR and H I or H₂ in that regime. In the future, we will quantify the effect the different subgrid physics modules and parameters have on the KS relation and its scatter.

To understand what gives rise to the different relations between Σ_{SFR} and the H I and H₂ gas content, it is informative to study where these gas phases are located in the phase-space diagram. We show this in Fig. 3 for galaxies with stellar masses $> 10^9 \text{ M}_{\odot}$ and $\text{SFR} > 0$. From the H I distribution we see the bimodality that

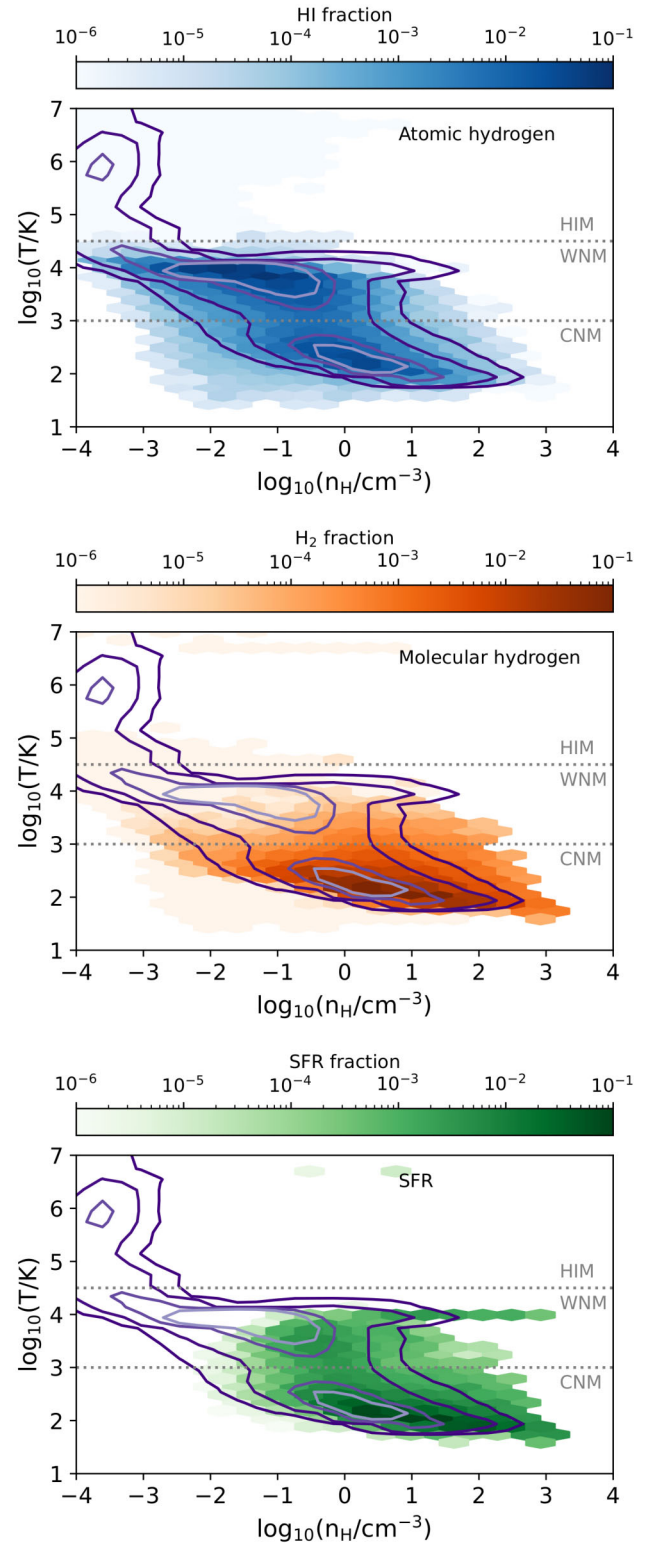


Figure 3. The distribution of all gas particles within 50 pkpc of the centre of mass of galaxies at $z = 0$ with $M_{\star} \geq 10^9 \text{ M}_{\odot}$ and $\text{SFR} > 0$ in the temperature-density plane. Contours enclose the regions where 99 per cent, 95 per cent, 68 per cent and 50 per cent of the particles are. The coloured hexbins show the H I (top), H₂ (middle) and SFR (bottom) contribution from each bin to the total H I, H₂ and SFR, respectively. The totals are obtained by summing these quantities of all the galaxies. The horizontal dotted lines mark the regions associated with the hot ionized (HIM), warm neutral (WNM) and cold neutral (CNM) media, as labelled, and as defined in equation (2).

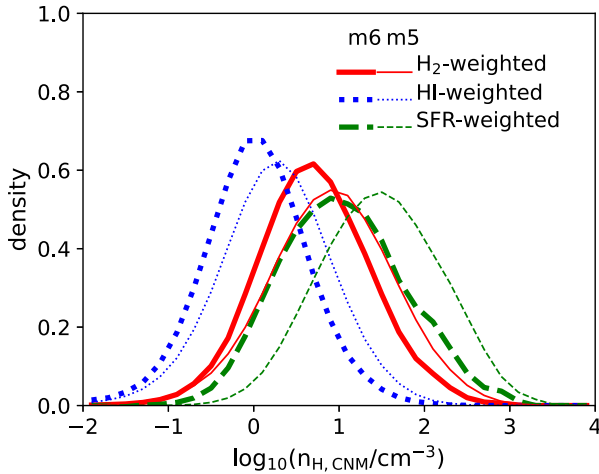


Figure 4. Probability density function of the gas density of gas particles in the CNM (as defined in equation 2), weighted by the H_2 , H I mass, or the SFR, as labelled. This is shown for $z = 0$ for gas particles in the L200 m6 (thick lines) and L025m5 (thin lines) simulations, that are in the CNM of galaxies at $z = 0$ with $M_\star \geq 10^9 M_\odot$ and $\text{SFR} > 0$.

arises from H I being primarily in the WNM and CNM (see equation 2) of the ISM. Most of the H I mass is however contributed by the WNM (≈ 64 per cent). For H_2 , we see that very little is in the WNM, with most of it being in the CNM (≈ 99 per cent). Note, however, that most of the gas mass in the CNM is atomic (≈ 70 per cent), while only ≈ 30 per cent is molecular. For the SFR, we see that ≈ 7 per cent is associated with the WNM and 93 per cent with the CNM.

Even though most of the CNM consists of atomic rather than molecular gas, the correlation of the SFR is stronger with H_2 than H I or total neutral gas (as seen by the smaller scatter of the H_2). The reason for this is that even within the CNM, the densities traced by H_2 and SFR are much higher than those traced by the H I . This is shown in Fig. 4 for the L200m6 and L025m5 simulations. Even in comparison to H_2 , the SFR is concentrated in higher density gas regardless of the resolution. Thus, the larger fraction of SFR associated with the high-density tail of the CNM, where the H_2 resides is the cause of the stronger correlation between Σ_{SFR} and Σ_{H_2} than between Σ_{SFR} and Σ_{HI} or $\Sigma_{\text{HI}+\text{H}_2}$. Comparing the m5 and m6 curves, we see that the SFR-weighted density distribution moves to higher values as the resolution increases. This is because the criterion for gravitational instability is evaluated at the resolution limit. The H I - and H_2 -weighted distributions also move to higher densities, but to a lesser extent than the SFR-weighted gas density distribution.

Fig. 3 shows a population of particles with $\text{SFR} > 0$ and $T \approx 10^4 \text{ K}$, which is neither associated with significant H_2 nor H I . These particles correspond to those that were evaluated to have non-zero SFR at the beginning of the time-step, but received energy from stellar feedback at the end of the time-step. This is a rare occurrence and these particles contribute only ≈ 0.7 per cent to the total instantaneous SFRs of galaxies in our sample. In any case, the star formation model in COLIBRE is such that these particles would not turn into stellar particles because they would be classified as non star-forming at the beginning of the next time-step.

3.2 KS relation: characterizing the scatter

3.2.1 Dependence on resolved galaxy properties

The KS relation predicted by COLIBRE and in the observations has a significant scatter. In this section we explore how the scatter is correlated with other resolved properties of galaxies, which were calculated using the particles tagged to belong to the same annuli in the ANNULI-FACE method.

Fig. 5 shows how the scatter in the KS relation at $z = 0$ for H I , H_2 , and total neutral hydrogen depends on the stellar surface density (top), gas metallicity (middle), and dust surface density (bottom). Overall, we find that in all gas phases, the scatter around the KS relation tends to decrease with increasing gas density, even for H I . However, the decrease is stronger for H_2 and total neutral gas.

For $\Sigma_\star$ (top panels in Fig. 5), we find that as $\Sigma_\star$ decreases, the local H_2 depletion time, computed as $\Sigma_{\text{H}_2}/\Sigma_{\text{SFR}}$ decreases, which is clear at $\Sigma_{\text{H}_2} \lesssim 10 M_\odot \text{ pc}^{-2}$. At higher H_2 surface densities, there is little dependence of the H_2 KS relation's scatter on $\Sigma_\star$. For H I , lower $\Sigma_\star$ are associated with lower H I -to-SFR conversion efficiencies at fixed Σ_{HI} , opposite to the trend displayed in the H_2 case. The trend between the KS relation scatter and $\Sigma_\star$ at fixed gas surface density appears stronger for H I than H_2 . We will come back to this point later. The total neutral gas appears similar to what is observed for the H I KS relation. This is not surprising as at $z = 0$ the overall gas content of galaxies tends to be dominated by H I rather than H_2 .

For the gas metallicity (i.e. metals not locked in dust; middle panels in Fig. 5), we see again opposite trends in H I and H_2 , the lower metallicity gas tends to be associated with a higher (lower) conversion efficiency between H_2 (H I) and stars. The dust surface density trends (lower panels in Fig. 5) resemble those of the gas metallicity: lower Σ_{dust} is associated with more (less) efficient conversion from H_2 (H I) into stars.

The physical drivers of the relations between the scatter of the KS relation and the gas metallicity and dust content are the same. As the gas becomes more metal and dust poor, it needs to reach higher densities for it to be cold and hence likely to form stars. In the absence of dust (or at least if it is scarce), the cold, dense gas will struggle to form H_2 . Hence, as the gas becomes more metal and dust-poor, less H_2 and more H I are associated with the same amount of SFR, leading to the appearance of a higher SFR conversion efficiency for H_2 , and a lower one for H I . M. E. Orr et al. (2018) found a similar tendency for low-metallicity gas to be associated with longer depletion time-scales for the total $\text{H I} + \text{H}_2$ KS relation in the FIRE simulations. They did not however explore the effect of gas metallicity on the H I and H_2 KS relation separately, partially because their model does not explicitly follow the formation of hydrogen species.

We show how the gas metallicity influences the H I and H_2 composition of the CNM in COLIBRE in Fig. 6. The H I mass fraction in the CNM saturates at ≈ 0.75 at low metallicity, as the rest of the gas is in helium and metals. The figure shows that the dependence of the H I fraction in the CNM on gas metallicity is qualitatively the same in the L200m6 and the L025m5 simulations, showing that the result is robust against resolution. We also show the typical gas density at which stars form in COLIBRE for the m6 and m5 resolution. At m6 resolution, the density consistently decreases with decreasing metallicity, even dropping below $n_{\text{H}} \approx 1 \text{ cm}^{-3}$ for the very metal-poor gas. This is, however, driven by resolution, as for the m5 run we see that star formation consistently happens at densities $> 10 \text{ cm}^{-3}$, even when H I overwhelmingly dominates the mass of CNM. We remind

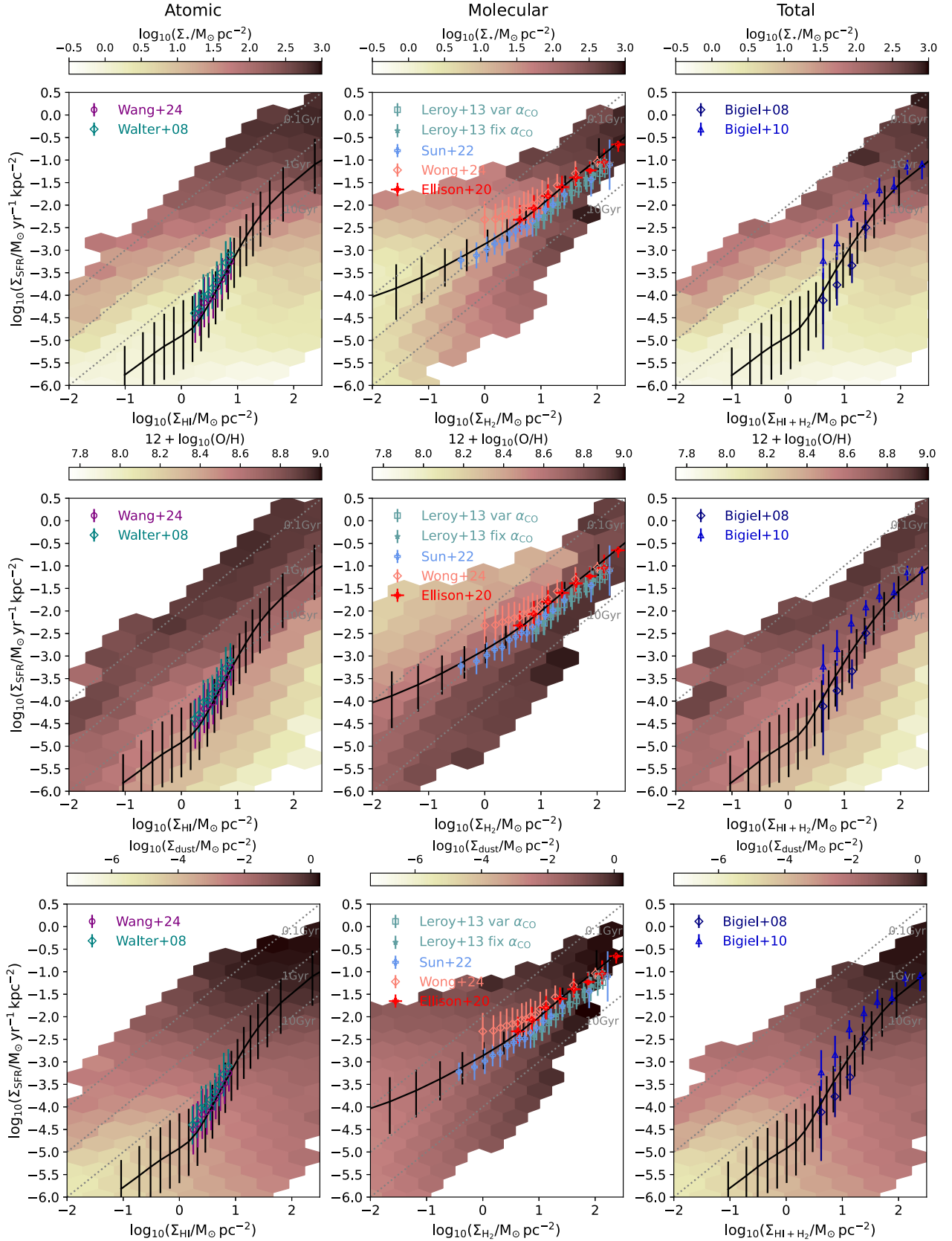


Figure 5. As Fig. 2 but with the coloured bins showing the median stellar surface density (Σ_* ; top panels), gas metallicity ($\log_{10}(\text{O}/\text{H})$; middle panels) and surface density of dust (Σ_{dust} , lower panels). Black lines with error bars in each panel show the median and 16th–84th percentile ranges, respectively. Only hexbins with ≥ 5 objects are included.

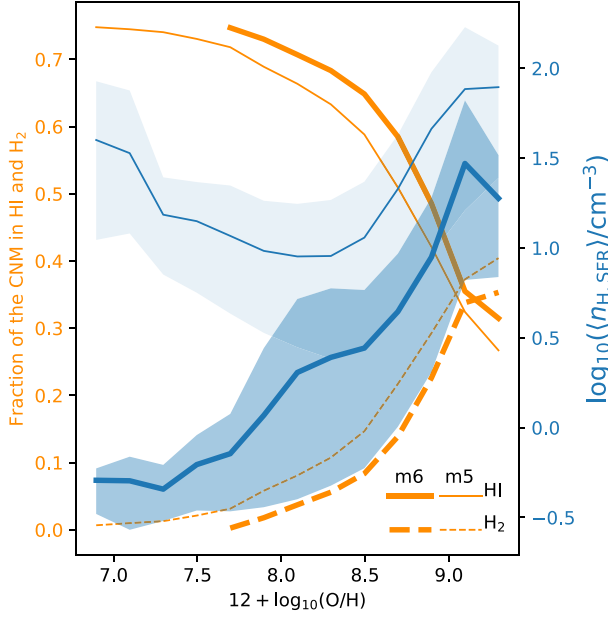


Figure 6. The fraction of the CNM in the form of H I (orange solid line) and H₂ (orange dotted line), and the SFR-weighted gas density (blue line with shaded region), as a function of gas metallicity. This is constructed with all the gas particles that are within 50 pkpc from the centres of galaxies with $M_* \geq 10^9 M_\odot$ and $\text{SFR} > 0$ at $z = 0$ in the L200m6 (thick lines) and L025m5 (thin lines) simulations. The H I fraction saturates at a CNM fraction of ≈ 0.75 , as the rest of the gas is in the form of helium and metals. For the gas density, we show the median and the 16th–84th percentile ranges.

the reader that this difference between resolutions is expected and intentional, as the instability criterion for star formation is evaluated at the resolution limit. At the Milky-Way metallicity of $12 + \log_{10}(\text{O}/\text{H}) \approx 8.75$, COLIBRE predicts the fraction of H I in the CNM to be ≈ 0.5 , which is similar to the value observed in the Milky-Way (N. M. McClure-Griffiths, S. Stanimirović & D. R. Rybarczyk 2023). This is indicative only, as we are not selecting galaxies to be of Milky-Way mass in Fig. 6. At lower metallicities, H.-J. Park et al. (2025) find that there is more H I in the CNM, which they interpret as the gas remaining H I as it gets colder due to the difficulty in forming H₂ in the absence of dust. This is in qualitative agreement with what Fig. 6 shows for COLIBRE.

Fig. 7 explores the strength of the correlations between the gas depletion time, measured in a spatially resolved manner, $\tau_{\text{gas}} \equiv \Sigma_{\text{gas}}/\Sigma_{\text{SFR}}$, and the resolved properties in galaxies (including those shown in Fig. 5). To show all the correlations together, we simply subtract from every property its corresponding median for the entire sample of Fig. 5. For example, for the stellar surface density, the x-axis shows $\log_{10}(\Sigma_*) - \langle \log_{10}(\Sigma_*) \rangle$.

For τ_{H_2} , we find that the strongest correlation is the negative correlation with specific Σ_{SFR} ($s\Sigma_{\text{SFR}}$; which is related to the correlation with Σ_*), and slightly weaker positive correlations appears with Σ_{dust} and local gas metallicity. The strong anti-correlation between τ_{H_2} and resolved sSFR means that the high sSFR is not just due to a higher concentration of cold gas but also that the cold gas is being converted into stars more efficiently. In the case of the gas metallicity and dust surface density, the presence of a moderately strong correlation with τ_{H_2} is not unexpected, as the accumulation of metals and dust directly affects the efficiency of gas cooling and H₂ formation, but the fact that

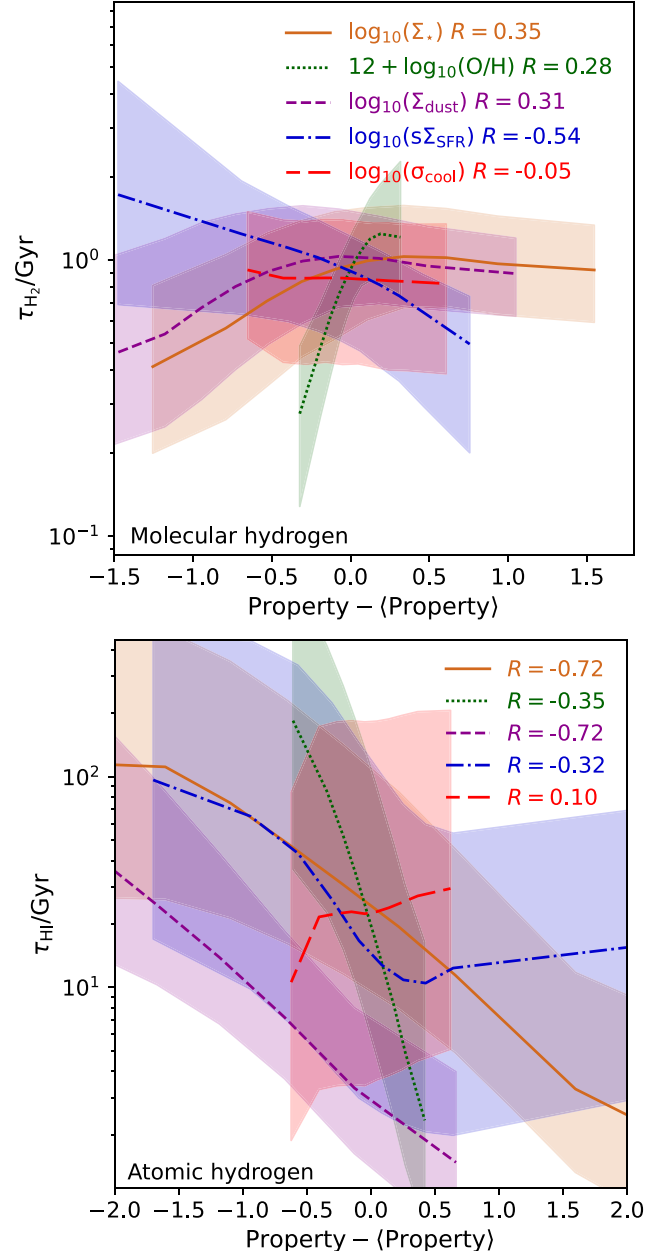


Figure 7. The correlation between the H₂ (top) and H I (bottom) depletion times, defined as $\Sigma_{\text{gas}}/\Sigma_{\text{SFR}}$, with gas being either H₂ or H I, and the stellar surface density, cold gas metallicity, surface density of dust, local sSFR and cool gas velocity dispersion (as defined in equation 3), as labelled. Lines and shaded regions show medians and 16th–84th percentile ranges, respectively, for galaxies at $z = 0$ with $M_* \geq 10^9 M_\odot$. To plot all the properties on the same x-axis scale, we subtract from every property its median. We also show the Spearman correlation coefficient, R , of each correlation, as labelled.

the correlations are positive is more surprising. In COLIBRE, dust acts both as a catalyst for the formation of H₂ and helps shield the gas. In this case, more of the hydrogen gas is in the form of H₂ and thus associated with a given SFR, giving the appearance of a longer H₂ depletion time. We test these predictions against observations in Section 3.2.3.

In the case of τ_{HI} (bottom panel in Fig. 7), we find that the strongest correlations are with the dust and stellar surface

densities (both with a Spearman correlation coefficient ≈ -0.75). The latter dependence has also been reported in observations; e.g. J. Wang et al. (2024) find $\tau_{\text{HI}} \propto \Sigma_{\star}^{-0.85 \pm 0.09} \Sigma_{\text{SFR}} / \Sigma_{\text{HI}}$. In COLIBRE, we find $\tau_{\text{HI}} \propto \Sigma_{\star}^{-0.72}$, in relatively agreement with these observations (see Section 3.2.3 for a more detailed comparison with J. Wang et al. 2024). The strong anti-correlation between τ_{HI} and Σ_{dust} is expected. The more dust there is, the easier it is for the hydrogen to transition from atomic to molecular, resulting in a relative lower fraction of star formation associated with atomic hydrogen. This leads to an apparent higher efficiency of H I to SFR conversion.

Interestingly, the cool gas ($T < 10^{4.5}$ K) one-dimensional velocity dispersion (equation 3) is only weakly correlated with τ_{HI} , while the correlation with τ_{H_2} is almost absent. J. Schaye et al. (2026) show that in COLIBRE, stars tend to form from cold gas with supersonic turbulent velocities. Our results show that, even though at the resolution limit the velocity dispersion (especially the turbulent component) play an important role in the star formation model, at kpc scales, star formation appears to be poorly correlated with gas velocity dispersion. We also tested different definitions of the gas velocity dispersion (weighting by H₂ or H I mass, and using different selections in distance from the mid-plane ranging from 3 to 10 kpc), and found that the weak correlation holds for different definitions.

The correlation with σ_{cool} manifests in a way that a higher velocity dispersion is associated with a longer H I depletion time. The star formation model in COLIBRE indirectly depends on the local gas velocity dispersion, since a higher value increases the Toomre instability parameter, making the gas less likely to collapse and be available for star formation (see F. S. J. Nobels et al. 2024 for a detailed discussion). The weaker correlation with σ_{cool} compared to some of the other local properties is in part due to σ_{cool} displaying less variations within and in-between galaxies, compared with other local properties. This is clear from the dynamic range covered by σ_{cool} which only extends over 1 dex, while $\Sigma_{\star}$, Σ_{dust} and $s\Sigma_{\text{SFR}}$ tend to display variations of $\gtrsim 1.5$ dex within our sample. Note, however, that the gas metallicity also varies over a small range compared to some of the other properties, while still displaying stronger correlations with both τ_{H_2} and τ_{HI} than what is obtained for σ_{cool} .

The total neutral hydrogen (H I + H₂) behaves very much like τ_{HI} , with the correlations only becoming a bit weaker than what is reported in the bottom panel of Fig. 7. This simply means that H I dominates the neutral gas reservoir in most galaxies at $z = 0$ in COLIBRE.

In Appendix A, we test the convergence of the correlations in Fig. 7 with numerical resolution. We find that all correlations are very well converged between the m6 and m5 resolution, with the exception being the correlation between τ_{H_2} and local gas metallicity. Although the correlation is similarly strong in the L025m5 compared to the L025m6 simulation (see Table A1), there is an offset in normalization. This is not surprising given that the stellar mass–gas metallicity relation is not well converged in COLIBRE at $z = 0$ (see discussion in J. Schaye et al. 2026).

3.2.2 Dependence on global galaxy properties

Fig. 8 presents the molecular gas KS relation at $z = 0$, separating the galaxy population by their global sSFR (top) and stellar mass (bottom). We show these properties because they are the ones that exhibited the strongest correlations with the scatter of the

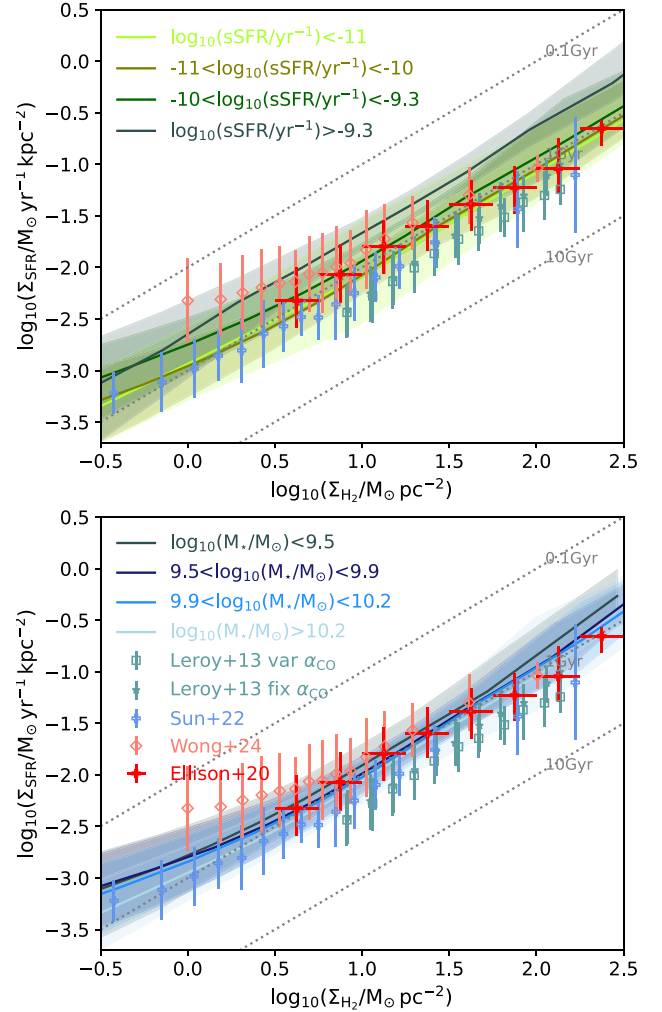


Figure 8. The H₂ KS relation at $z = 0$ for galaxies selected based on their sSFR (top panels) and stellar mass (bottom panel). The median and 16th–84th percentiles are shown with solid lines and shaded regions, respectively. All galaxies included here have $M_{\star} \geq 10^9 M_{\odot}$ and $\text{SFR} > 0$.

KS relation. In addition to the trends in Fig. 8, we also explored a possible dependence of the KS relation scatter on global gas metallicity, stellar age, galaxy type (whether the galaxies are defined as central or satellites) and galaxy morphology, and found only weak or no correlations.

The top panel of Fig. 8 shows a strong dependence of the molecular KS relation’s normalization on the global sSFR of the galaxy, in a way that galaxies with higher sSFR have shorter molecular gas depletion times. COLIBRE galaxies can have high sSFRs from a combination of increased gas reservoirs and a more efficient conversion of cold gas into stars. We quantify this effect by measuring the Spearman correlation coefficient R , between the resolved molecular depletion time and the global sSFR, at fixed global molecular and stellar masses. We find $R \approx -0.6$ with negligible p-values, indicating strong anticorrelations. The dependence across bins of stellar and molecular gas masses is similar and takes the shape $\tau_{\text{H}_2} \propto s\text{SFR}^{\beta}$, with $\beta \approx [-0.8, -0.9]$. Thus, COLIBRE predicts that changes in the global sSFR are accompanied by changes in the local conversion efficiency of molecular gas to stars.

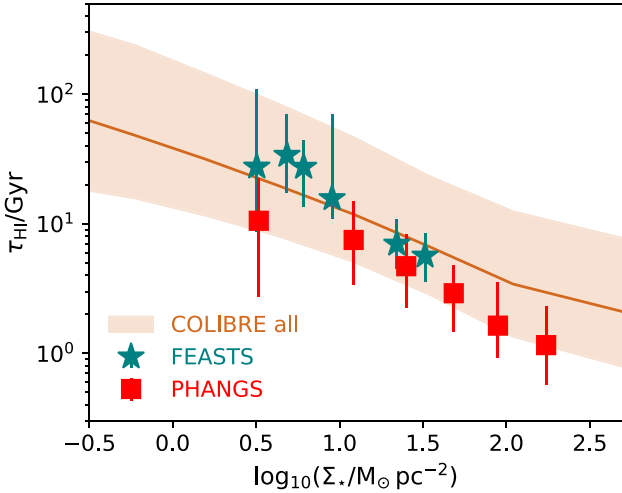


Figure 9. The dependence of the (resolved) H I depletion time, τ_{HI} on the local stellar surface density, Σ_* , for COLIBRE galaxies at $z = 0$ with $M_* \geq 10^9 M_\odot$ compared with the FEASTS (J. Wang et al. 2024) and PHANGS (J. Sun et al. 2022) surveys. The lines and symbols show the median of COLIBRE and the observations, respectively, while the shaded region and error bars show the 16th – 84th percentile ranges. This threshold matches well the minimum Σ_{HI} value in FEASTS and PHANGS.

The bottom panel of Fig. 8 shows a weak dependence on stellar mass, so that lower stellar mass galaxies have slightly elevated Σ_{SFR} at fixed Σ_{H_2} , with the differences being more clearly seen at $\Sigma_{\text{H}_2} \gtrsim 10 M_\odot \text{pc}^{-2}$. We find that the dependence on stellar mass in the bottom panel of Fig. 8 is driven by the sSFR dependence, as the sSFR has a negative dependence on stellar mass for main sequence galaxies (see fig. 16 in J. Schaye et al. 2026).

3.2.3 Comparison with observations

In this section, we present a series of comparisons between the predicted dependence of the scatter in the KS relation on galaxy properties, starting with the atomic gas and then moving to the molecular gas. Because most observational constraints are dedicated to the molecular KS relation, this section is devoted primarily to that.

Y. Shi et al. (2011) showed that there is a strong dependence of the scatter in the atomic gas KS relation on the stellar surface density, suggesting that it could be fundamentally driven by the dynamical equilibrium pressure. We compare the predicted relation between the local H I depletion time ($\equiv \Sigma_{\text{HI}}/\Sigma_{\text{SFR}}$) and the local stellar surface density with observations in Fig. 9. We show observations from the FEASTS (J. Wang et al. 2024) and PHANGS (J. Sun et al. 2022) surveys. FEASTS comprises 17 spiral galaxies with stellar masses $10^{9.7} - 10^{11} M_\odot$, while PHANGS contains 70 nearby galaxies with stellar masses between 10^9 and $10^{11} M_\odot$. In both observational sets, the relation was computed using radial annuli with size ≈ 1 kpc, similar to our method. The agreement between COLIBRE and the observations is excellent, even when we did not attempt to match the exact distribution of galaxy properties of the two surveys and limited to sampling a similar stellar mass range. This relation is very well converged with resolution, as shown in Fig. A3, so the agreement extends to the other COLIBRE simulations.

As shown in Fig. 7, Σ_* is one of the most significant drivers of the scatter in the H I KS relation in COLIBRE, and the same

appears to be the case in the observations. It is important to note that FEASTS appears to be slightly elevated compared to PHANGS, while COLIBRE lies in between. This most likely comes from the diffuse H I FEASTS is sensitive to (due to the data coming from the single dish Five-hundred-meter Aperture Spherical Telescope), but that PHANGS misses (as it primarily employs the interferometric Very Large Array). According to J. Wang et al. (2024), diffuse H I can account for up to 40 per cent of the H I mass, which directly translates into increasing τ_{HI} in a direct proportion. In COLIBRE, the local dynamical equilibrium pressure does not enter the star formation model, implying that the $\tau_{\text{HI}} - \Sigma_*$ correlation does not necessarily directly emerge from that parameter.

The scatter of the molecular gas KS relation in galaxies in the local Universe has been well studied with the ALMaQUEST (L. Lin et al. 2019), EDGE-CALIFA (T. Wong et al. 2024), and PHANGS (J. Sun et al. 2022) surveys. Here, we compare our predictions in detail with ALMaQUEST and PHANGS. The ALMaQUEST sample consists of 46 galaxies with stellar masses $10^{9.7} - 10^{11.5} M_\odot$, with 25 of those being on the main sequence and 21 being below the main sequence (referred to as ‘Green Valley’ galaxies). Via a combination of MaNGA and ALMA CO(2–1) observations, they characterize the SFR, molecular gas, stellar surface density and gas metallicity of individual spaxels in galaxies, with spaxels having sizes of 0.3 kpc. The PHANGS sample with H₂ profile measurements consists of 61 galaxies (stellar mass range is as indicated above), and here we use the data provided in annuli of 1 kpc width from J. Sun et al. (2022). Using the same sSFR criterion as L. Lin et al. (2019), the PHANGS catalogue consists of 19 main sequence and 45 Green Valley galaxies, thus PHANGS samples lower sSFR galaxies than ALMaQUEST on average. Another important difference between PHANGS and ALMaQUEST is that the SFRs of the regions in PHANGS are characterized using a combination of H α , FUV, NUV, and MIR, while in ALMaQUEST they come only from H α .

Although PHANGS extends to lower stellar masses compared with ALMaQUEST, most galaxies (54/61) have masses $\gtrsim 10^{9.8} M_\odot$. Thus, to compare with observations, we select galaxies in COLIBRE at $z = 0$ that have stellar masses $> 10^{9.7} M_\odot$ and $\text{SFR} > 0$. We choose a slightly lower stellar mass limit than the minimum mass of the ALMaQUEST sample to allow for potential systematic uncertainties in the derivation of stellar masses of the order of those inferred at $z = 0$ (S. Bellstedt & A. S. G. Robotham 2025). To perform the comparison with observations, we define $\Delta \Sigma_{\text{SFR}} \equiv \log_{10}(\Sigma_{\text{SFR}}) - \langle \log_{10}(\Sigma_{\text{SFR}}(\Sigma_{\text{H}_2})) \rangle$, where Σ_{SFR} corresponds to the value of an annulus (or spaxel in the case of ALMaQUEST) and $\langle \log_{10}(\Sigma_{\text{SFR}}(\Sigma_{\text{H}_2})) \rangle$ corresponds to the median of all annuli (or spaxels) within a bin of Σ_{H_2} . We calculate this quantity for both simulation and observations, which helps us to focus on how the scatter of the relation correlates with properties of interest.

We start by comparing the resolved dependence of the scatter of the molecular KS relation on the sSFR in Fig. 10. Here, we take all $z = 0$ galaxies in COLIBRE with stellar masses $> 10^{9.7} M_\odot$, and simply bin annuli by their local sSFR $\equiv \Sigma_{\text{SFR}}/\Sigma_*$. COLIBRE predicts a very strong dependence of the scatter in the molecular KS relation on the local sSFR that agrees remarkably well with the observed dependence in ALMaQUEST at $\text{sSFR} > 10^{-10} \text{yr}^{-1}$, which better samples the high sSFR region. At lower sSFR values, we see that ALMaQUEST displays a lower normalization of the KS relation than predicted by COLIBRE, while there is

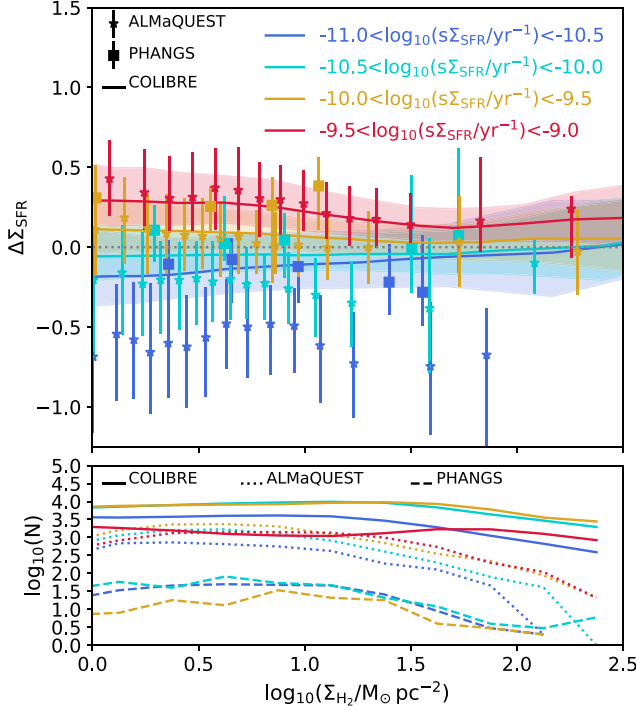


Figure 10. *Top panel:* Deviations from the median molecular gas KS relation ($\log_{10}(\Sigma_{\text{SFR}}) - \langle \log_{10}(\Sigma_{\text{SFR}}) \rangle$) at $z = 0$ as a function of H_2 surface density, separating regions by their local sSFR ($s\Sigma_{\text{SFR}} \equiv \Sigma_{\text{SFR}}/\Sigma_*$, for observations (symbols) and COLIBRE (lines). Symbols and lines show medians, while error bars and shaded regions show the 16th – 84th percentile ranges. For ALMaQUEST (L. Lin et al. 2019), we include all galaxies and spaxels; for PHANGS (J. Sun et al. 2022), we include all galaxies and annuli; while for COLIBRE we include $z = 0$ galaxies with $M_* > 10^{9.7} M_\odot$ and $\text{SFR} > 0$. Note that no PHANGS data are presented for the highest $s\Sigma_{\text{SFR}}$ bin, as there were not enough statistics to produce reliable measurements for the medians and scatter. The agreement between COLIBRE and the observations is remarkable. At low values of sSFR, COLIBRE agrees better with PHANGS. *Bottom panel:* histogram of Σ_{H_2} for each of the samples of the top panel. The observational and simulated samples follow similar distributions, with most regions having $\log_{10}(\Sigma_{\text{H}_2}/M_\odot \text{ pc}^{-2}) \lesssim 1.4$. This is also the regime in which the differences between sSFR samples are clearest in both observations and simulations.

excellent agreement with PHANGS. This may be related to ALMaQUEST using $\text{H}\alpha$ to derive SFRs, which is known to systematically underestimate SFRs compared to UV/IR tracers at low SFRs (J. C. Lee et al. 2009). PHANGS on the other hand, by using a combination of tracers, circumvents this limitation.

We remind the reader that beyond the local sSFR, the scatter of the molecular KS relation is correlated with the local gas metallicity (Fig. 7). We test this prediction with ALMaQUEST only, as the data for the metallicity profiles for PHANGS were not available (see J. Sun et al. 2022 for details). Fig. 11 shows $\Delta \Sigma_{\text{SFR}}$ as a function of Σ_{H_2} for samples selected by local gas metallicity. To build Fig. 11, we removed all spaxels in the ALMaQUEST sample that have too low signal-to-noise ratio ($\text{S/N} < 3$) to measure a gas metallicity, which removes 17 per cent of the spaxels. The removed spaxels tend to be those with long H_2 depletion times. Thus, the resulting sample tends to be biased towards short depletion times and high Σ_{H_2} (see the bottom panel of Fig. 11) compared to what is seen in the bottom panel of Fig. 10. Despite these differences, we see a similar trend in COLIBRE and

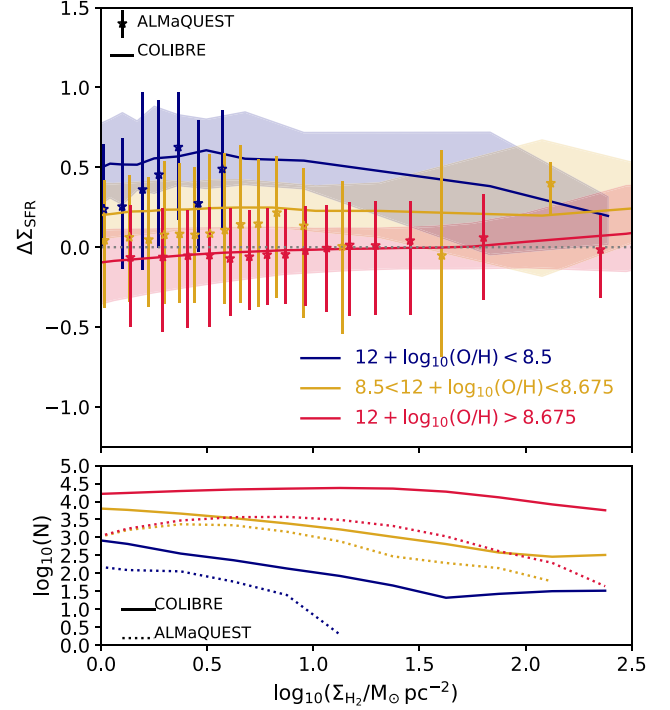


Figure 11. As Fig. 10 but for regions binned by the local gas metallicity, $12 + \log_{10}(\text{O}/\text{H})$, and comparing only with ALMaQUEST. The sampling of the metallicity bins is similar in both observations and simulations, with a similar distribution of $\log_{10}(\Sigma_{\text{H}_2}/M_\odot \text{ pc}^{-2})$ in each metallicity bin.

ALMaQUEST: $\Delta \Sigma_{\text{SFR}}$ decreases as the gas metallicity increases. This is particularly clear in the $0.25 \lesssim \log_{10}(\Sigma_{\text{H}_2}/M_\odot \text{ pc}^{-2}) \lesssim 0.75$ range that is best sampled by the observations.

We stress that from the ALMaQUEST sample alone, it is not clear whether gas metallicity is more weakly or strongly correlated with τ_{H_2} than other properties are. This could in part be due to the large fraction of spaxels we need to remove from the sample due to the low signal-to-noise ratio, but also to the limited stellar mass range of the sample, which is skewed towards massive galaxies. The mass–metallicity relation is steep at $M_* \lesssim 10^{10} M_\odot$ but flattens at higher stellar masses (e.g. C. A. Tremonti et al. 2004). Hence, the overall metallicity range sampled by ALMaQUEST is rather small, with 90 per cent of the spaxels having $8.57 < 12 + \log_{10}(\text{O}/\text{H}) < 8.8$. Future surveys sampling a larger dynamic range in gas metallicity would be needed to test the relative importance of different properties in setting the scatter of the molecular KS relation.

For the comparison above, we used all spaxels in ALMaQUEST and all annuli in PHANGS regardless of whether galaxies were classified as star-forming or green valley, and only investigated the dependence of the scatter in the molecular KS relation on resolved properties. L. Lin et al. (2022) argue that global galaxy properties also impact the resolved KS relation, and find that green-valley galaxies in ALMaQUEST have a lower global molecular-to-star formation conversion efficiency than star-forming galaxies. In Fig. 12, we compare with those results by splitting our COLIBRE sample into star-forming and green valley galaxies, in such a way that the median sSFR of each sample resembles that of ALMaQUEST and PHANGS. We find that a selection $s\Sigma_{\text{SFR}} > 10^{-9.75} \text{ yr}^{-1}$ in COLIBRE provides a median sSFR $\approx 10^{-9.64} \text{ yr}^{-1}$ while the median for the star-forming

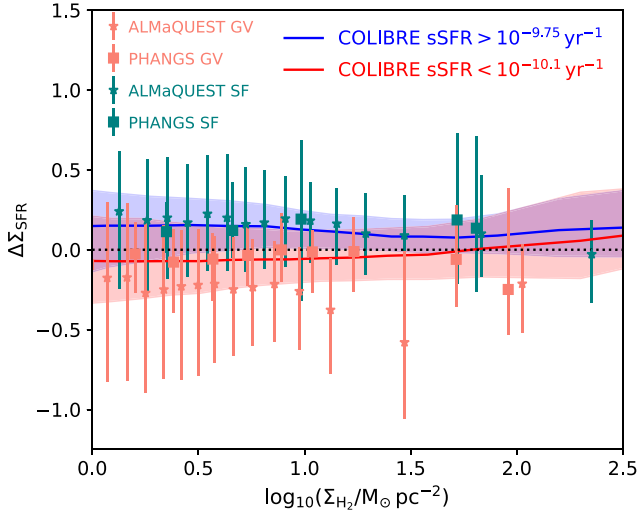


Figure 12. As the top panel of Fig. 10, but for star-forming and green-valley galaxies separately, in COLIBRE (lines) and observations (symbols), as labelled. This is shown for galaxies at $z = 0$. Symbols and lines show medians, while error bars and shaded regions show the 16th–84th percentile ranges. To select the two samples in COLIBRE, we apply thresholds in $sSFR$, as labelled, which give median $sSFR$ similar to the median $sSFR$ in ALMaQUEST and PHANGS for the two samples. The agreement between COLIBRE and the observations is very good, particularly when comparing with PHANGS.

samples in ALMaQUEST and PHANGS is $sSFR \approx 10^{-9.65} \text{ yr}^{-1}$ and $\approx 10^{-9.7} \text{ yr}^{-1}$, respectively. We select green-valley galaxies in COLIBRE to have $sSFR < 10^{-10.1} \text{ yr}^{-1}$, which produces a sample with a median $sSFR \approx 10^{-10.65} \text{ yr}^{-1}$, while the green-valley ALMaQUEST and PHANGS samples have medians of $sSFR \approx 10^{-10.57} \text{ yr}^{-1}$ and $\approx 10^{-10.45} \text{ yr}^{-1}$, respectively. We remind the reader that in COLIBRE, annuli are selected to have $SFR > 0$. We find excellent agreement between the molecular KS relation in COLIBRE with both sets of observations for the SF galaxies, while for green-valley galaxies, the agreement is better with PHANGS. The latter resembles the findings of Fig. 10, where the $H\alpha$ -derived SFRs at low $sSFR$ in ALMaQUEST appear to underestimate the SFR values compared to the multiwavelength approach of PHANGS.

The comparisons presented in Figs 2, 10, 11, and 12 show that COLIBRE predicts a resolved KS relation at $z = 0$ in different gas phases in excellent agreement with the observations. This is not only the case for the average relation, but also for the dependence of the KS relation’s scatter on local and global galaxy properties. This is a remarkable success, as the KS relation was not used to calibrate the free parameters in COLIBRE.

3.3 KS relation: galaxy-to-galaxy variations

We have explored the median KS relation and how the scatter correlates with resolved and global galaxy properties. In this section, we explore how much variation there is in the KS relation from galaxy to galaxy at $z = 0$. To quantify this, we fit the relationship between Σ_{SFR} and Σ_{gas} , with gas here being HI, H_2 or the sum of the two, using the function:

$$\Sigma_{SFR} = A_0 \Sigma_{gas}^{N_{KS}} \quad (4)$$

In this equation, Σ_{SFR} and Σ_{gas} take the units $M_\odot \text{ yr}^{-1} \text{ kpc}^{-2}$ and $M_\odot \text{ pc}^{-2}$, respectively. The fits for COLIBRE are performed using only gas with a surface density $\geq 0.1 M_\odot \text{ pc}^{-2}$. This is done to avoid the flattening of the KS relation at lower surface densities due to resolution limitations (see Fig. A1) and to focus on regions that are directly comparable with observations. We do not see large variations in the individual fits as long as they are performed with a minimum surface density in the range $0.1\text{--}1 M_\odot \text{ pc}^{-2}$. In addition, we only fit galaxies with a KS relation sampled with ≥ 5 rings. This means that we do not fit some of the low SFR galaxies in COLIBRE. The example galaxy shown at the bottom-right panel of Fig. B1 has a well-defined fit for its HI KS relation, but not for the H_2 KS relation. We find that the requirement above removes 37 per cent of the galaxies with $M_\star \geq 10^9 M_\odot$ and $SFR > 0$ at $z = 0$. The removed galaxies tend to have lower stellar masses and $sSFR$ s (with medians $\approx 10^{9.48} M_\odot$ and $\approx 10^{-10.17} \text{ yr}^{-1}$, respectively) than the ones that are kept for the analysis (with medians $\approx 10^{9.73} M_\odot$ and $\approx 10^{-10.02} \text{ yr}^{-1}$, respectively). We tested different minimum numbers of bins between 3–7 and found our results to be insensitive to that change.

Fig. 13 shows the correlations of N_{KS} and A_0 with stellar mass. For HI (blue), we find that the median slope of the entire population is $N_{KS} \approx 2.15$, but the value of individual galaxies strongly anticorrelates with the stellar mass of the galaxy. Low-mass galaxies tend to have higher values of N_{KS} , which is also seen for fits to the total gas (not shown here). The scatter around the median N_{KS} tends to increase with decreasing stellar mass for fits based on HI or total gas. For HI, the zero-point of the KS fits is positively correlated with stellar mass. For the molecular gas, we see that the slope and the zero-point are close to constant, with values ≈ 1 and ≈ -2.8 , respectively (in the units used in equation 4). We also find virtually no difference between the different COLIBRE resolutions.

The number of fits to the observed KS relation of individual galaxies is limited to a few dozens, even in the local Universe. We show in Fig. 13 observational results from different surveys, which agree on a close to constant KS molecular gas relation. We highlight that the observations of M. J. Jiménez-Donaire et al. (2023) correspond to Virgo cluster galaxies, with some of them being gas-deficient and clearly affected by their environment. Nevertheless, they follow a very similar relation. It is notable that the observations suggest a similar scatter in the slope and zero-point of the molecular KS relation, as predicted by COLIBRE. The overall trends seen, including a potential weak dependence on stellar mass, also agree well with COLIBRE’s predictions.

For HI, we again see very good agreement between the m5 and m6 resolutions. We note that the difference in A_0 for the HI KS relation between the L200m6 and L025m5 runs at around $M_\star \approx 10^{9.5} M_\odot$ is driven by the different box sizes, as the differences become negligible when we compare the L025m5 simulation with the L025m6 simulation instead of L200m6 (not shown here). It is informative to show the L025m5 run along with the L200m6 run as the observational fits to the HI KS relation are biased towards dwarf galaxies, which are HI dominated, and better resolved in the m5 run. To compare COLIBRE with observations, we compiled data from F. Bigiel et al. (2008), S. Roychowdhury et al. (2009), J. Wang et al. (2024). S. Roychowdhury et al. (2009) did not report stellar masses for their sample of dwarf galaxies, and hence we computed our own stellar masses from IRAC 3.6 μm photometry. We collected the IRAC 3.6 μm photometry from the

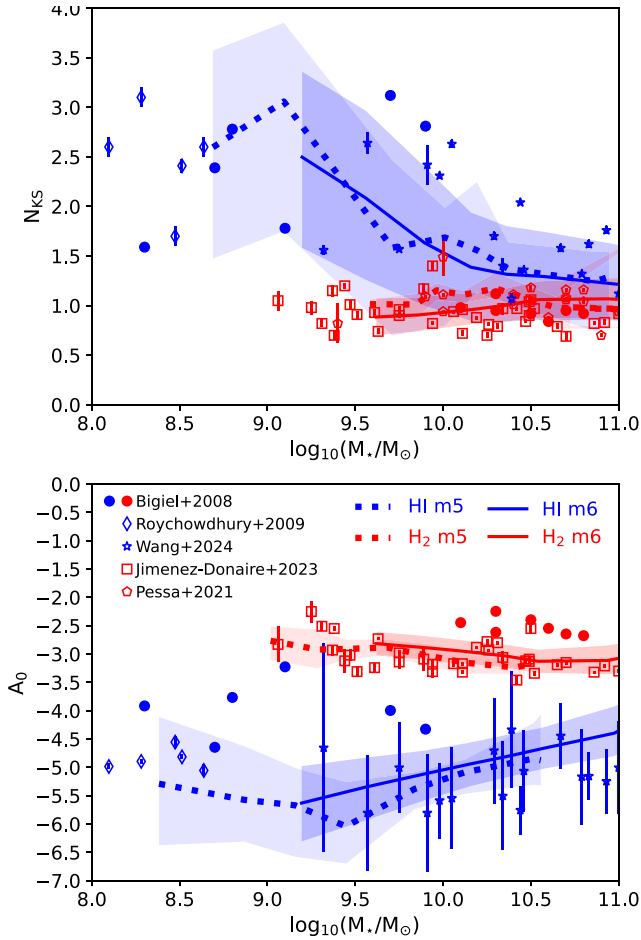


Figure 13. The slope (top) and zero-point (bottom) of the KS relation fits for individual $z = 0$ galaxies as a function of stellar mass in COLIBRE. We show results for the L200m6 (solid lines) and L025m5 (dotted lines) simulations, as labelled in the bottom panel. The equation being fitted is equation (4). We show the slope and zero-point of the fits for the case of fitting only H I (blue) and H₂ (red). Lines with shaded regions show the median and the 16th–84th percentile ranges, respectively. Observational estimates from F. Bigiel et al. (2008), S. Roychowdhury et al. (2009), I. Pessa et al. (2021), M. J. Jiménez-Donaire et al. (2023), J. Wang et al. (2024) are shown as symbols, as labelled in the bottom panel, using the same colour-code of COLIBRE.

NASA/IPAC Extragalactic Database for the S. Roychowdhury et al. (2009) sample, and adopted a mass-to-light ratio corresponding to a G. Chabrier (2003) IMF following A. K. Leroy et al. (2008). Although the scatter is large, the observations are consistent with a dependence of the slope and zero-point of the H I KS relation on stellar mass that is similar to what COLIBRE predicts. Because the predicted scatter is so large, much larger observational samples would be required to validate or refute our predictions.

4 REDSHIFT EVOLUTION OF THE KENNICUTT–SCHMIDT RELATION

4.1 The evolving KS relation from $z = 0$ to $z = 8$

Fig. 14 shows the evolution of the KS relation in H I, H₂ and total neutral hydrogen, from $z = 8$ to $z = 0$. COLIBRE predicts a monotonic evolution of the H₂ KS relation with time, so that the

normalization of the relation consistently decreases with time, maintaining a power-law index $N_{KS} \approx 1$ at all redshifts (at least at $\log_{10}(\Sigma_{H_2}/M_\odot \text{ pc}^{-2}) \gtrsim 1$). The zero-point of the H₂ KS relation evolves very strongly with redshift, decreasing by a factor of ≈ 20 from $z = 8$ to $z = 0$.

The case of H I and total neutral hydrogen is more complex: at low gas surface densities $\lesssim 100 M_\odot \text{ pc}^{-2}$, the normalization of the KS relation decreases mildly with increasing redshift from $z = 0$ to $z = 8$ (by ≈ 0.8 dex for H I and 0.4 dex for H I + H₂), while also becoming slightly steeper. This evolution reverses at a gas surface density $\approx 100 M_\odot \text{ pc}^{-2}$, resembling the evolution of the H₂ KS at higher surface densities.

Interestingly, we do not see a strong evolution in the scatter of the H₂ KS relation, which is inconsistent with the findings reported in other simulations, such as NEWHORIZON. K. Kraljic et al. (2024) find that the scatter in the H₂ KS relation in NEWHORIZON decreases by ≈ 0.25 dex from $z = 4$ to $z = 0.25$. There are several differences between our work and theirs; for example, they include all galaxies with stellar masses $\geq 10^7 M_\odot$, while here we adopt a mass threshold that is 100 times higher. Another important difference is that NEWHORIZON does not directly model the formation of H₂, and hence it is assumed that thresholds in temperature and density are able to exclusively select H₂ gas. We demonstrated in Figs 3 and 6 that this is not the case and, in fact, as the metallicity decreases, more of the cold gas can remain atomic. These differences make it hard to pinpoint the driver of the discrepancies but they certainly serve as a motivation to constrain these relations observationally. On the other hand, X. Shen et al. (2026) presented an analysis of the KS relation for total neutral gas at $z > 3$ for Thesan-zoom simulations measured at 1 kpc-scales. They found a steep relationship at high surface densities, $\Sigma_{H I + H_2} \gtrsim 10^{1.2} M_\odot \text{ pc}^{-2}$, which agrees well with ours. At lower surface densities, their simulation predicts much higher Σ_{SFR} compared to COLIBRE, but this is a regime where Thesan-zoom is not well sampled given the small number of galaxies and the overall smaller number of low density regions with $SFR > 0$. The steep neutral KS relation is interpreted by X. Shen et al. (2026) as the result of turbulent energy balance in the ISM maintained by stellar feedback. The importance of turbulence in regulating star formation is also predicted by COLIBRE, as reported by J. Schaye et al. (2026).

The strong evolution of the zero-point of the H₂ KS relation in COLIBRE is in part due to the evolution of the average gas metallicity of galaxies with $M_* \geq 10^9 M_\odot$. Fig. 15 shows the fraction of CNM in H I and H₂ as a function of the gas metallicity from $z = 0$ to $z = 5$ (similar to what Fig. 6 shows), and the evolution of the gas metallicity when weighted by different properties. First, we see that, at all redshifts, there is a clear trend of an increasing amount of the CNM being H I as the metallicity decreases. However, this relation is redshift dependent, with a decreasing fraction of the CNM in H I with increasing redshift at fixed gas metallicity. The corresponding simultaneously decreasing metallicity with redshift, leads to an overall effect of less H₂ and more H I being associated with a given amount of SFR with increasing redshift. This gives the impression that the molecular gas is more efficient at forming stars at high redshift, but this is primarily due to H₂ becoming a poorer tracer of star formation at high redshift, at least on kpc scales. We remind the reader that the star formation model in COLIBRE assumes a universal star formation efficiency per free fall time of 1 per cent and a universal gravitational instability criterion to select star-forming gas (equation 1), and hence there is no evolution built

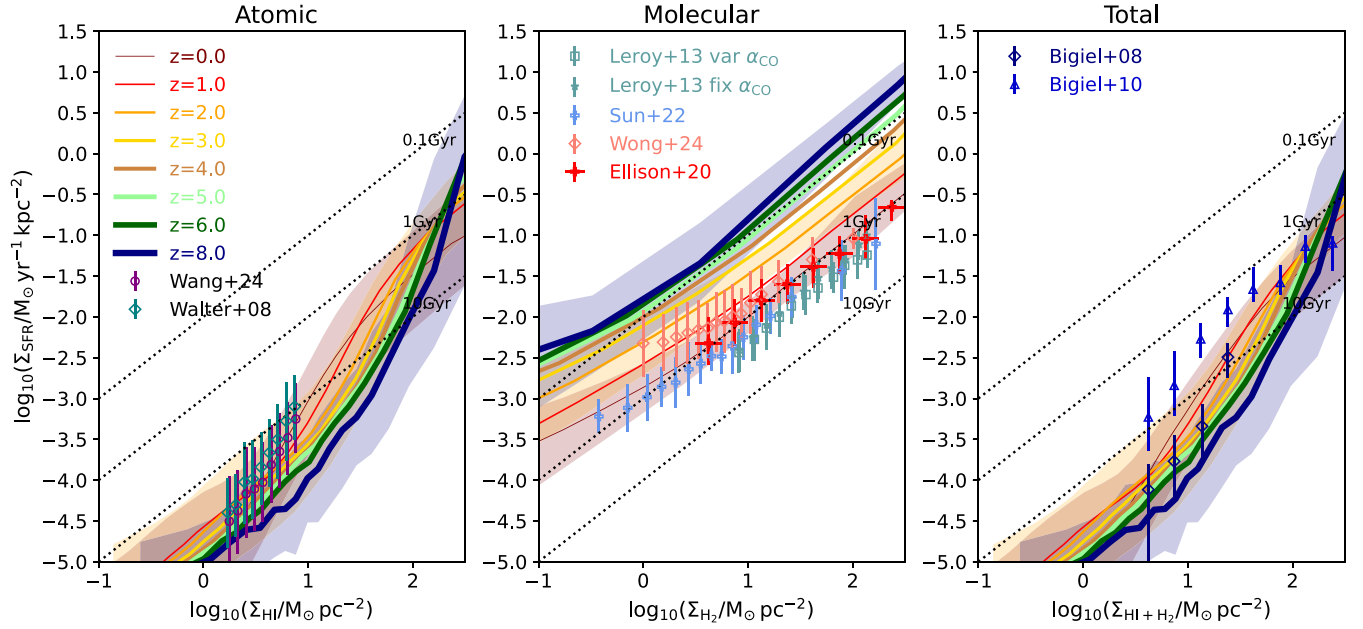


Figure 14. As Fig. 2 but showing the evolution of the median KS relation for galaxies with $M_{\star} \geq 10^9 M_{\odot}$ and $\text{SFR} > 0$ from $z = 8$ to $z = 0$. We show the 16th–84th percentiles only for $z = 0, 2$, and 8 for clarity. Redshifts are as labelled in the left panel. The observational data shown in all panels corresponds to $z \approx 0$.

into the model. Thus, the strong evolution in the KS relation predicted by COLIBRE, particularly for H_2 , is driven by the processes that govern the formation of H_2 and determine whether the gas is unstable and what the density distribution of the unstable gas is.

Fig. 15 shows that the gas metallicity weighted by the H I mass is always lower than the values weighted by SFR or H_2 , with the differences becoming larger at higher redshifts. This is not surprising, as H I traces lower density gas than H_2 and SFR, which is typically located in the outskirts of galaxies, with most galaxies in COLIBRE having negative metallicity gradients (see Serrano Rodríguez et al. in preparation).

4.2 Comparison with observations of the resolved and global KS relation

In Fig. 16, we present comparisons with observations of the KS relation at $z > 0$. There are only very limited samples of galaxies at intermediate and high redshifts for which one can measure a resolved KS relation, and they correspond primarily to strongly lensed galaxies (e.g. J. A. Hodge et al. 2015; C. E. Sharon et al. 2013, 2019; D. Nagy et al. 2023). Note that at $z = 1$ the observations correspond to 6 galaxies (4 coming from J. Freundlich et al. 2013 and 2 from D. Nagy et al. 2023). At $z = 2$, there are observations for four galaxies, one is a submillimetre highly star-forming galaxy (C. E. Sharon et al. 2013); two are a pair of hyper-luminous IR galaxies (J. S. Gómez et al. 2025) separated by ≈ 85 kpc, one of which is mildly magnified; and a main-sequence galaxy (C. E. Sharon et al. 2019). At $z = 4$, the observations correspond to five galaxies (four from M. Béthermin et al. 2023 and 1 from J. A. Hodge et al. 2015). In all cases, a constant (Milky-Way or starburst galaxy) CO-to- H_2 conversion factor was used. The conclusions we can draw from such a small sample are limited, but can serve as a reference to compare COLIBRE predictions with. We applied

corrections to the SFR to convert them from a Salpeter or Kroupa IMF to a Chabrier IMF when necessary, and remove the helium contributions from the molecular gas if quantities reported include it.

At $z = 1$, we find that the global measurements of D. Nagy et al. (2023) agree well with the predicted relation for COLIBRE, but the individual regions (which sample a 800 pc scale) display a very large scatter, larger than the typical variations in COLIBRE. The resolved relations measured for the Cosmic snake and the sample of J. Freundlich et al. (2013) agree well with COLIBRE, but the individual regions in A521 tend to scatter down towards longer H_2 depletion times. At $z = 2$, we see that COLIBRE’s median relation agrees better with the main sequence galaxy in C. E. Sharon et al. (2019), while the highly starbursting galaxies of C. E. Sharon et al. (2013) and J. S. Gómez et al. (2025) lie (slightly) above the predicted relation. This shows that even within the small observational samples, there are hints that the normalization of the KS relation varies with a galaxy’s sSFR, in agreement with our predictions. At $z = 4$, we find that the observations sample Σ_{H_2} to values higher than what is typically reached by COLIBRE at m6 resolution, but the extrapolation of the predicted relation agrees well with what J. A. Hodge et al. (2015) report and is potentially above what M. Béthermin et al. (2023) measure. Larger samples are required to start to systematically probe the evolution of the resolved KS relation. This will become possible in the near future, for example with the extended Northern Extended Millimetre Array (NOEMA).

Unlike spatially resolved observations, integrated values for individual galaxies have been obtained by the thousands (L. J. Tacconi et al. 2020). From these observations, it has been well established that the H_2 depletion time decreases with increasing redshift. This is shown for COLIBRE and a compilation of observations in Fig. 17. The observations come from a range of H_2 tracers: CO, sub-millimetre or millimetre emission, or FIR SED.

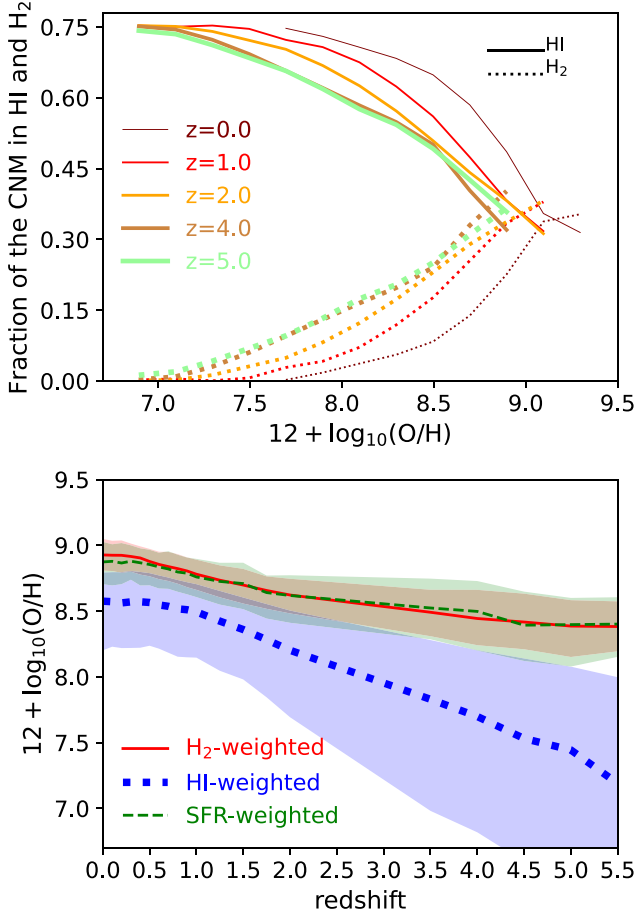


Figure 15. *Top panel:* The fraction of the CNM in HI (solid lines) and H₂ (dotted lines) as a function of gas metallicity, from $z = 0$ to $z = 5$ (see equation 2 for a definition of the CNM). We apply the same stellar and SFR selection of Fig. 6. Note that we only include a subset of the redshifts shown in Fig. 14 to avoid overlapping lines. *Bottom panel:* The redshift evolution of the median gas metallicity (solid lines), and the 25th – 75th percentile range (shaded regions), of galaxies with $M_* \geq 10^9 M_\odot$, when weighted by the H₂ mass (solid line), HI mass (dotted line), or SFR (dashed line).

Even though they are all subject to different systematic effects, L. J. Tacconi et al. (2020) argue that they produce consistent H₂ scaling relations. Each circle represents an individual observed galaxy, while the stars show the median for the observations. The latter are computed with a minimum of 10 galaxies. We show with lines the predictions of COLIBRE, which are in remarkable agreement with the observations. However, we remind the reader that for COLIBRE we show individual 1 kpc rings rather than global measurements for galaxies, so the comparison is not one-to-one. Another important caveat is that we show regions from all galaxies with $M_* \geq 10^9 M_\odot$ and $\text{SFR} > 0$, while the observations are biased towards more massive galaxies with higher SFRs. Below we explore the effect of sSFR.

The scatter in the observations is large and driven by physical parameters, such as the sSFR of galaxies, which we show in Fig. 8 can systematically shift the H₂ depletion time. We demonstrate this at different redshifts in Fig. 18, where we plot the resolved H₂ depletion time as a function of a galaxy’s sSFR in five redshift bins. At all redshifts, there is a strong dependence of the local τ_{H_2}

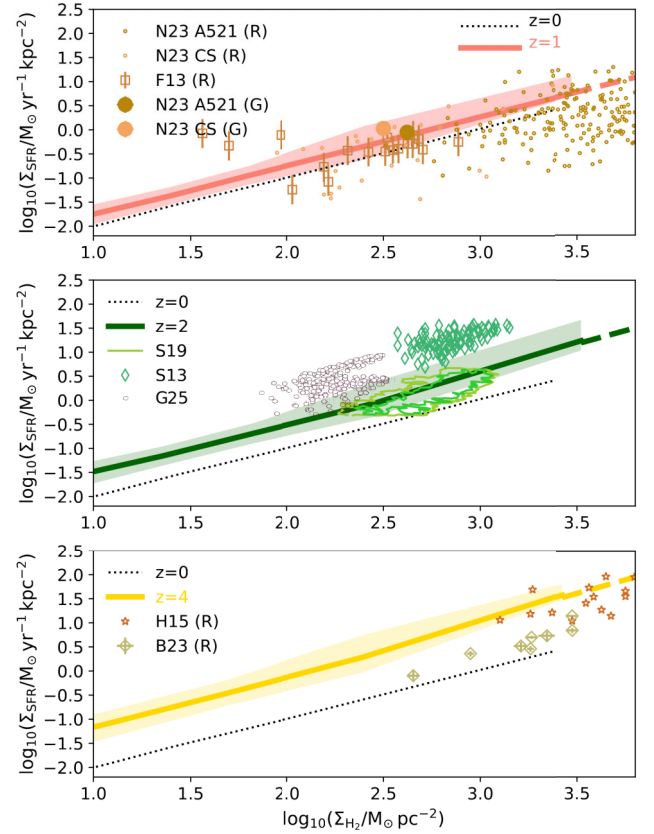


Figure 16. The predicted KS relation at $z = 1$ (top), $z = 2$ (middle) and $z = 4$ (bottom) for galaxies with $M_* \geq 10^9 M_\odot$. We compare with a compilation of observations of the (resolved) KS relation in individual galaxies. At $z = 1$, we show the resolved (R) and global (G) measurements of A521 and the Cosmic Snake from D. Nagy et al. (2023) (N+23). For the resolved data in N+23, we use the data provided in 800 pc apertures. We also show the resolved data from four massive star-forming galaxies at $z \approx 1.2$ from J. Freundlich et al. (2013). At $z = 2$ we compare with four galaxies; two from C. E. Sharon et al. (2013) (S13; a submillimetre galaxy at $z = 2.6$) and C. E. Sharon et al. (2019) (S19; a main sequence galaxy at $z = 2.26$); and a pair of two hyperluminous IR galaxies from J. S. Gómez et al. (2025) (G25). For S13, we show the individual ≈ 2 kpc pixels; for S19 we show the contours tracing one and three pixels, where each pixel is of ≈ 2 kpc size (see C. E. Sharon et al. 2019 for details); while for G25 we show individual ≈ 2 kpc pixels for each galaxy. At $z = 4$, we show the results from three massive, star-forming galaxies from the ALPINE survey (M. Béthermin et al. 2023) (B23) at $z \approx 4.5$, and one submillimetre galaxy at $z \approx 4$ J. A. Hodge et al. (2015) (H15). For COLIBRE, we show as dashed lines the extrapolation at high surface densities of the linear fit to the median relation (solid lines) measured at $1 \leq \log_{10}(\Sigma_{\text{H}_2}/M_\odot \text{ pc}^{-2}) \leq 3$.

on the global sSFR, with the zero-point of the relation evolving strongly with redshift in COLIBRE. We also compute medians for the observations in redshift bins where there is enough data to have at least 10 points per bin. The observations agree strikingly well with our predictions, noting though that the observational samples are not stellar-mass complete, which our simulated sample is by construction. We also note that at the highest sSFR, the observations prefer slightly longer H₂ depletion times than what COLIBRE predicts at kpc scales. In the future, we will compare global measurements to perform a fairer comparison with these global observations.

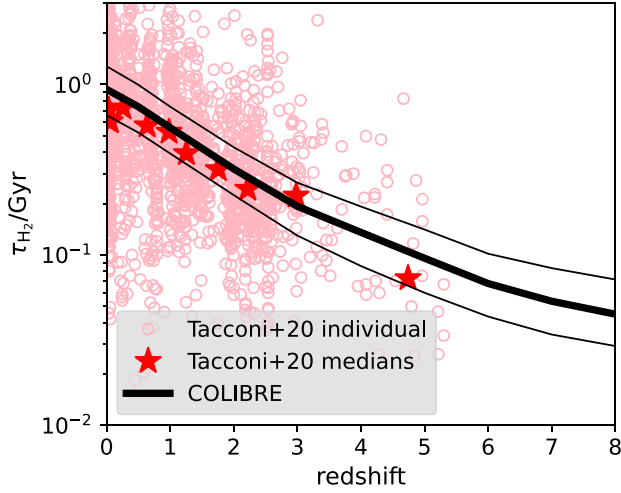


Figure 17. The evolution of the H_2 depletion time compared with the compilation of observations presented in L. J. Tacconi et al. (2020). Small circles correspond to individual galaxies, while the stars correspond to the medians computed from the observations. The thick and thin solid lines correspond to the median and 16th–84th percentile ranges of all galaxies in COLIBRE with $M_* \geq 10^9 M_\odot$ and $SFR > 0$. To compute τ_{H_2} , we only considered 1 kpc annuli with $\Sigma_{H_2} > 1 M_\odot pc^{-2}$, which correspond to the gas density regime where the L200m6 run is well converged (see Fig. A1) and comfortably includes the regimes from which observational results are computed.

4.3 The evolving dependence of the scatter in the KS relation with local physical properties

Finally, in Fig. 19 we explore how some of the strongest correlations between τ_{H_2} and τ_{HI} with other local properties presented in Fig. 7 for $z = 0$ evolve with time. For τ_{H_2} , we find that the correlation with local gas metallicity remains strong with increasing redshift, with the amplitude of the correlation being close to redshift invariant. However, the correlation displays a saturation at the highest metallicities at any one redshift, with a reversal of the correlation at the high-metallicity end. We confirmed that the τ_{H_2} -gas metallicity correlation does not change when we increase the resolution, nor when we adopt instead the hybrid AGN feedback model (using the L025m5 and L050m6 Hybrid runs listed in Table 1). When using the L025m5 run, we find that the relation holds and the saturation manifests itself at the same τ_{H_2} values as in the L200m6 run. Interestingly, this saturation happens at $\approx 0.1/H(z)$, where $H(z)$ is Hubble’s constant at redshift z . The latter corresponds to the free-fall time of a sphere of density $\sim 200 \rho_{crit}$, where ρ_{crit} is the critical density of the universe. Realistically, star formation cannot proceed any slower than that timescale, as $200 \rho_{crit}$ corresponds to the minimum gas density for gravitational collapse in a CDM universe. The redshift invariance of the τ_{H_2} -gas metallicity correlation implies that the physical mechanism behind it remains the same throughout cosmic time: as gas metallicity decreases, more of the star-forming gas is in the form of HI, and hence the amount of H_2 associated with a fixed SFR decreases. This was shown in Fig. 15.

The bottom-left panel of Fig. 19 shows that the correlation between τ_{H_2} and $s\Sigma_{SFR}$ remains strong even up to $z = 8$, but with the normalization of the relation decreasing strongly with redshift. This mirrors the strong evolution we see in the relation

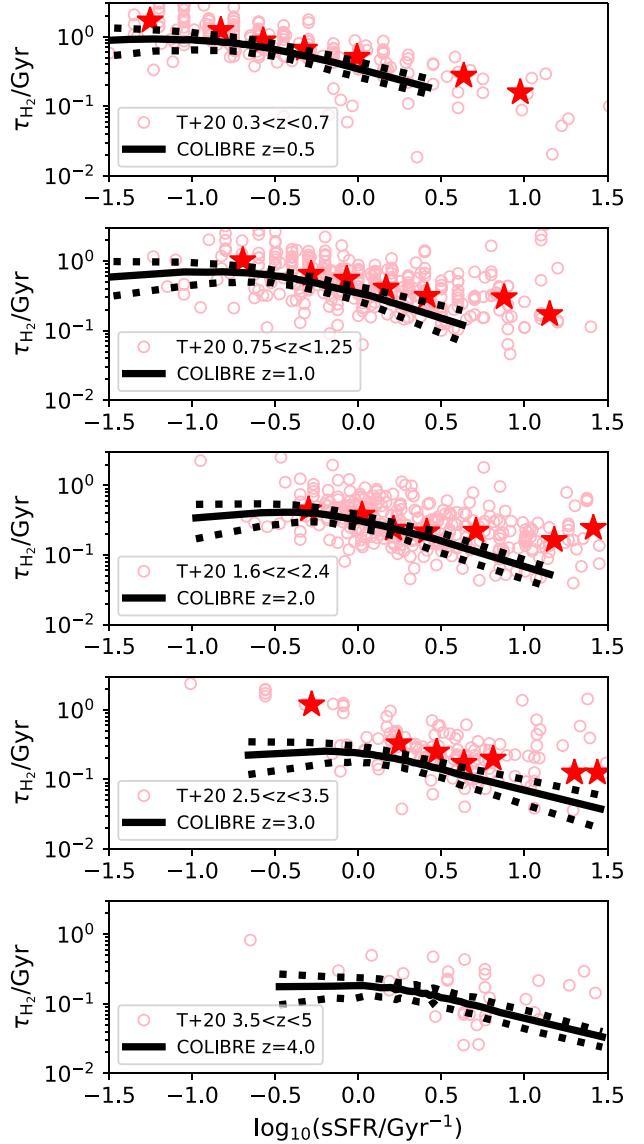


Figure 18. The H_2 depletion time as a function of $sSFR$ for galaxies with $M_* \geq 10^9 M_\odot$, at different redshift bins from $z \approx 0.5$ to $z \approx 5$, compared with a compilation of observations presented in L. J. Tacconi et al. (2020). The solid and dashed lines in each panel show the median and 16th – 84th percentile ranges, respectively, for galaxies in COLIBRE. Symbols are as in Fig. 17, and their corresponding redshift range are labelled in each panel. We do not show medians for the observations for the bottom panel due to the small number of galaxies.

between τ_{H_2} and the global $sSFR$ of galaxies in Fig. 18. Ignoring the saturation, the close to redshift invariance of the τ_{H_2} -metallicity relation suggests a more fundamental physical connection than that between τ_{H_2} and $s\Sigma_{SFR}$, even if statistically the latter may be stronger in some regimes (as measured by the Spearman correlation coefficient).

In the top-right panel of Fig. 19, we show that the correlation between τ_{HI} and metallicity evolves very strongly with redshift, with τ_{HI} decreasing by two orders of magnitude from $z = 0$ to $z = 8$ at fixed local gas metallicity. The bottom-right panel of Fig. 19 shows that the relation between τ_{HI} and Σ_* is already in place at $z = 8$ in COLIBRE, and evolves more mildly with redshift than the relation with local gas metallicity, with

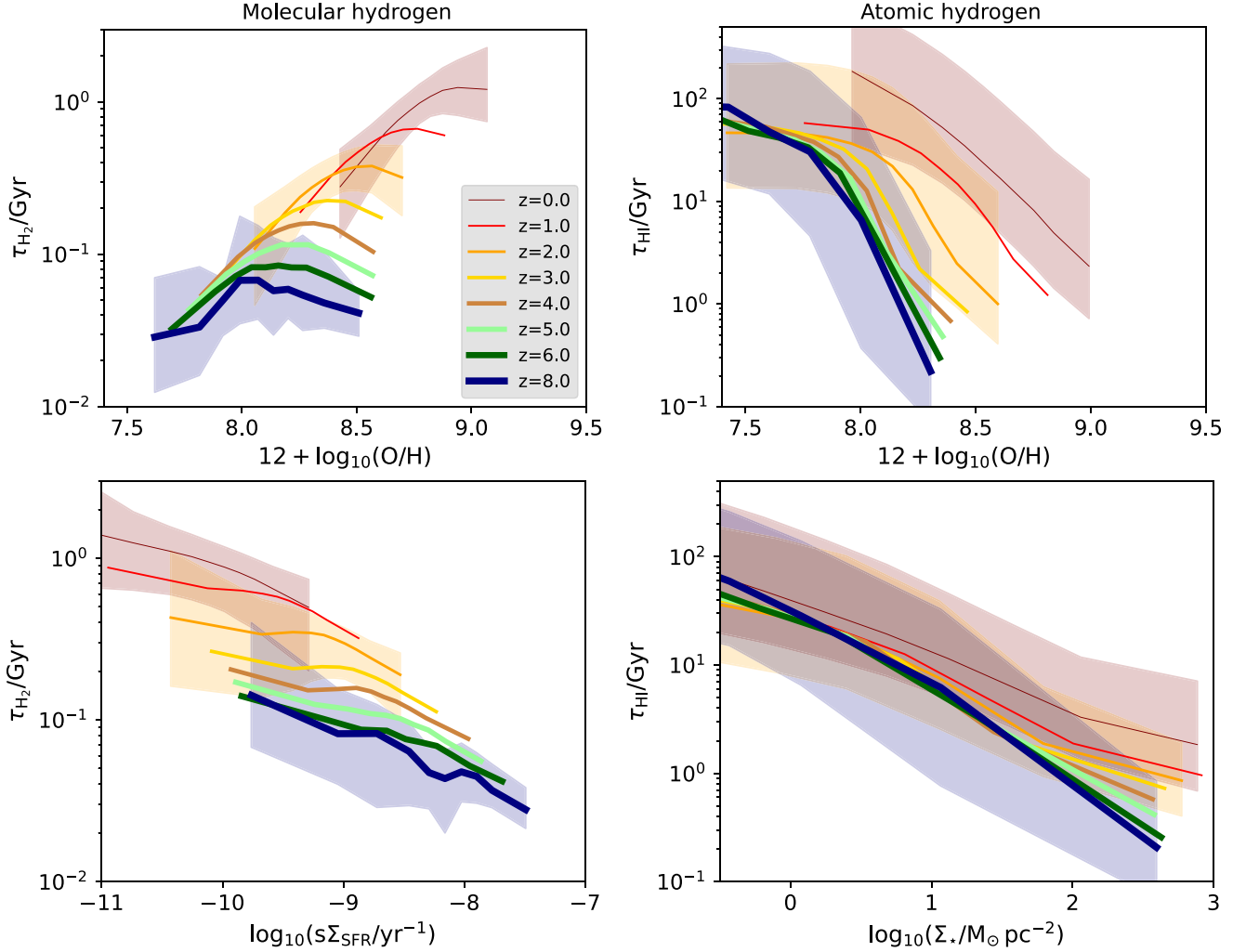


Figure 19. Evolution of some of the strongest resolved relations found in COLIBRE, between local galaxy properties and the kpc-scale H_2 (left) and H I (right) depletion timescales, for galaxies with $M_\star \geq 10^9 M_\odot$. For H_2 , we show the evolution of the relation between the resolved τ_{H_2} and the local gas metallicity (top-left panel) and sSFR surface density (bottom-left panel), between $z = 0$ and 8. Solid lines with shaded regions represent the medians with 16th–84th percentile ranges (only shown for $z = 0, 2$, and 8), respectively, with the colours indicating different redshifts, as labelled in the top-left panel. For H I , we show the relation with local gas metallicity (top-right panel) and the stellar surface density (bottom-right panel).

τ_{HI} decreasing by a factor of ≈ 3 from $z = 0$ to $z = 8$ at fixed $\Sigma_\star$. For H I , there is usually larger scatter in all the relations we explored compared with H_2 , which is natural, as a significant fraction ($\gtrsim 50$ per cent) of the H I is not in the CNM, but instead in the WNM, which is less directly related to star formation.

Overall, our results show that the kpc-scale KS relation is already in place by $z = 8$ in COLIBRE. This is true for both the H I and H_2 KS relations. This contrasts with the predictions presented in K. Kraljic et al. (2024) for NEWHORIZON, in which they found that the KS relation is well established only at $z \lesssim 2-3$. However, as we mentioned above, there are important differences between our sample and theirs in terms of the stellar mass of galaxies covered, while also having several differences in the physical models. In the future, we want to investigate the impact on the KS relation of switching on and off various processes in COLIBRE, including, but not limited to, early stellar feedback, the non-equilibrium chemistry and the chemical diffusion model.

5 DISCUSSION AND CONCLUSIONS

We used the COLIBRE cosmological hydrodynamical simulations of galaxy formation to investigate the H I , H_2 , and total neutral hydrogen KS relations and the scatter of these relations for galaxies with $M_\star \geq 10^9 M_\odot$, over $0 \leq z \leq 8$. COLIBRE features on-the-fly non-equilibrium chemistry for H and He , with metals also included, in equilibrium but with non-equilibrium electron abundances, detailed radiative cooling down to ≈ 10 K, and live dust that impacts gas cooling and the formation of H_2 . This detailed modelling enables direct predictions for the H I , H_2 , and total neutral H KS relations.

We measure the KS relation using circular rings of fixed width in galaxies oriented face-on. We find, however, that modifying the method to measure the KS relation has only minimal effects on our results (Appendix A3). Our main findings are summarized below:

- (i) COLIBRE accurately reproduces the observed H I , H_2 , and total gas KS relations at $z = 0$ (Fig. 2), including their scatter,

without having been explicitly tuned to these data. COLIBRE's star formation model assumes an underlying constant gas-to-star formation efficiency per free-fall time of 1 per cent, with the gas required to be gravitationally unstable at the resolution limit to form stars. This model, coupled with the sophisticated ISM model in COLIBRE, naturally captures differences in the KS relation amplitude and scatter between H I and H₂, with H I showing longer depletion times and larger scatter than H₂, consistent with the H I's association with more diffuse gas phases (Figs 3, 4, and 6). The kpc-average KS relations are very well converged in COLIBRE, even when the numerical resolution is varied by 2 orders of magnitude.

(ii) COLIBRE predicts the scatter in the KS relation to correlate strongly with local properties such as the stellar surface density, gas metallicity, and dust surface density (Figs 5 and 7). At kpc scales, we find that lower metallicity/dust content regions form stars efficiently before significant H₂ formation, leading to apparently shorter H₂ and longer H I depletion times. The strongest predictors of H₂ depletion time are the local sSFR and gas metallicity, while H I depletion time correlates most strongly with the stellar surface density and dust content (Fig. 7).

(iii) COLIBRE agrees remarkably well with spatially resolved observations at $z = 0$. In particular, we compared with ALMaQUEST, PHANGS, and J. Wang et al. (2024) and found broad agreement in both the KS relation itself and its dependence on resolved properties such as sSFR and metallicity (Figs 9–11). COLIBRE also reproduces the lower KS amplitude in green-valley galaxies relative to main-sequence systems, matching ALMaQUEST and PHANGS (Fig. 12). We also explore KS fits performed on a galaxy-by-galaxy basis and find that COLIBRE predicts variations in the best-fitting parameters that resemble what observations indicate for H I and H₂ (Fig. 4), including a dependence of the best-fitting parameters on stellar mass for H I.

(iv) While the global stellar mass plays a minor role in the normalization of the KS relation (Fig. 8), the galaxy-wide sSFR has a stronger influence: high-sSFR galaxies have both larger cold gas reservoirs and higher star-formation efficiencies. We find that the enhanced efficiency is not purely due to gas content but also linked to higher conversion rates of gas into stars.

(v) The kpc-scale H I, H₂, and H I + H₂ KS relations are already in place by $z = 8$ in COLIBRE. At even higher redshifts, there are too few galaxies that meet our selection ($M_* > 10^9 M_\odot$ and $SFR > 0$) to clearly establish the presence of a KS relation. The H₂ depletion time decreases by a factor of ≈ 20 from $z = 0$ to $z = 8$ (Fig. 14), largely driven by the lower local gas metallicities at higher redshift, leading to more H I-dominated CNM from which stars form (Fig. 15). The monotonic evolution of the H₂ KS relation contrasts with that of the H I, which shows a more complex evolution, with a transition at around $\Sigma_{\text{HI}} \approx 100 M_\odot \text{ pc}^{-2}$. At lower gas surface densities, the amplitude of the H I KS relation decreases mildly with redshift, while at higher densities, the trend reverses.

(vi) The overall lower gas metallicities at high redshifts predicted by COLIBRE lead to higher star formation efficiencies per unit H₂ in early galaxies, which is consistent with recent high-redshift constraints (Figs 16–18). The lower gas metallicity is associated with a CNM that is H I dominated, which once averaged on kpc scales, looks like having an efficient H₂ to SFR conversion. COLIBRE predicts a strong correlation between τ_{H_2} and sSFR, which agrees very well with observational constraints coming from integrated galaxy properties.

Our results demonstrate that the KS relation at kiloparsec scales for the H I and H₂ gas is an emergent outcome of the ISM physics (cooling, dust grains, metal enrichment, stellar feedback, and non-equilibrium chemistry) and the underlying universal star formation efficiency included in the COLIBRE model. This is the case not only for the median KS relation, but also for the dependence of the scatter on local and global galaxy properties, and its very strong redshift evolution. The agreement with spatially resolved observations of the KS relations is a significant success of the COLIBRE model because none of the model's free parameters were adjusted to match these observations.

In the future we want to use variations of the COLIBRE physics, e.g. assuming chemical equilibrium, turning off the live dust model, varying the parameters of the star formation model and the stellar feedback, to isolate the physical mechanisms setting the amplitude and scatter of the KS relation at different cosmic epochs. This will also allow us to understand whether the KS relation can be used to study the interplay between galaxy-scale feedback effects and star formation.

One crucial limitation is the small observational samples of galaxies for which a resolved KS relation has been measured, and the ad hoc galaxy selections. This complicates the testing of our predictions. Upcoming observations with ALMA of hundreds of local Universe galaxies² will allow the H₂ KS relation to be measured at kiloparsec scales for stellar-mass selected samples, and a robust characterization of the scatter and its dependence on other local and global properties. Similarly, at higher redshifts, the combination of JWST and NOEMA/ALMA will allow increasing the sample of galaxies for which a resolved KS relation can be measured to dozens.

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We acknowledge the use of SOAP (R. McGibbon et al. 2025) in our analysis, and thank their developers for the effort put in the code and documentation.

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²<https://www.kilogas.space/>

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DATA AVAILABILITY

The data supporting the plots within this article are available on reasonable request to the corresponding author. The COLIBRE simulation data will eventually be made publicly available, although we note that the data volume (several petabytes) may prevent us from simply placing the raw data on a server. In the meantime, people interested in using the simulations are encouraged to contact the corresponding author. A public version of the SWIFT code (M. Schaller et al. 2024) is available at <http://www.swiftsim.com>. The COLIBRE modules implemented in SWIFT will be made publicly available after the public release of the simulation data.

REFERENCES

- Abbott T. M. C. et al., 2022, *Phys. Rev. D*, 105, 023520
- Agertz O., Kravtsov A. V., Leitner S. N., Gnedin N. Y., 2013, *ApJ*, 770, 25
- Agertz O. et al., 2021, *MNRAS*, 503, 5826
- Ali-Haïmoud Y., Bird S., 2013, *MNRAS*, 428, 3375
- Aoyama S., Hou K.-C., Shimizu I., Hirashita H., Todoroki K., Choi J.-H., Nagamine K., 2017, *MNRAS*, 466, 105
- Bahé Y. M. et al., 2022, *MNRAS*, 516, 167
- Bellstedt S., Robotham A. S. G., 2025, *MNRAS*, 540, 2703
- Benítez-Llambay A. et al., 2026, *MNRAS*, 546, stag268
- Béthermin M. et al., 2023, *A&A*, 680, L8
- Bigiel F., Leroy A., Walter F., Brinks E., de Blok W. J. G., Madore B., Thornley M. D., 2008, *AJ*, 136, 2846
- Bigiel F., Leroy A., Walter F., Blitz L., Brinks E., de Blok W. J. G., Madore B., 2010, *AJ*, 140, 1194
- Bigiel F. et al., 2016, *ApJ*, 822, L26
- Blitz L., Rosolowsky E., 2006, *ApJ*, 650, 933
- Booth C. M., Schaye J., 2009, *MNRAS*, 398, 53
- Borrow J., Schaller M., Bower R. G., Schaye J., 2022, *MNRAS*, 511, 2367
- Chabrier G., 2003, *PASP*, 115, 763
- Chaikin E., Schaye J., Schaller M., Benítez-Llambay A., Nobels F. S. J., Ploekinger S., 2023, *MNRAS*, 523, 3709
- Chaikin E. et al., 2026, *MNRAS*, 548, stag300
- Chandro-Gómez Á. et al., 2025, *MNRAS*, 539, 776
- Correa C. A. et al., 2026, *MNRAS*, 548, stag645
- Dalla Vecchia C., Schaye J., 2012, *MNRAS*, 426, 140
- Davé R., Anglés-Alcázar D., Narayanan D., Li Q., Rafieeantsoa M. H., Appleby S., 2019, *MNRAS*, 486, 2827
- Davé R., Crain R. A., Stevens A. R. H., Narayanan D., Saintonge A., Catinella B., Cortese L., 2020, *MNRAS*, 497, 146
- Dehnen W., 2014, *Comput. Astrophys. Cosmol.*, 1, 1
- Dell’Aglì F., García-Hernández D. A., Schneider R., Ventura P., La Franca F., Valiante R., Marini E., Di Criscienzo M., 2017, *MNRAS*, 467, 4431
- Dey B. et al., 2019, *MNRAS*, 488, 1926
- Diemer B., Sparre M., Abramson L. E., Torrey P., 2017, *ApJ*, 839, 26
- Ellison S. L. et al., 2020, *MNRAS*, 493, L39
- Faucher-Giguère C.-A., 2020, *MNRAS*, 493, 1614
- Federrath C., 2015, *MNRAS*, 450, 4035
- Federrath C., Klessen R. S., 2012, *ApJ*, 761, 156
- Feldmann R. et al., 2023, *MNRAS*, 522, 3831
- Forouhar Moreno V. J., Helly J., McGibbon R., Schaye J., Schaller M., Han J., Kugel R., 2025, *MNRAS*, 543, 1339
- Freundlich J. et al., 2013, *A&A*, 553, A130
- Gallagher M. J. et al., 2018, *ApJ*, 858, 90
- Glover S. C. O., Clark P. C., 2012, *MNRAS*, 421, 9
- Gómez J. S., Messias H., Nagar N. M., Orellana-González G., Ivison R. J., van der Werf P., 2025, *A&A*, 704, A178
- Gonzalez-Perez V., Lacey C. G., Baugh C. M., Lagos C. D. P., Helly J., Campbell D. J. R., Mitchell P. D., 2014, *MNRAS*, 439, 264
- Granato G. L. et al., 2021, *MNRAS*, 503, 511
- Greif T. H., Glover S. C. O., Bromm V., Klessen R. S., 2009, *MNRAS*, 392, 1381
- Han J., Jing Y. P., Wang H., Wang W., 2012, *MNRAS*, 427, 2437
- Han J., Cole S., Frenk C. S., Benítez-Llambay A., Helly J., 2018, *MNRAS*, 474, 604
- Hirashita H., Voshchinnikov N. V., 2014, *MNRAS*, 437, 1636
- Hodge J. A., Riechers D., Decarli R., Walter F., Carilli C. L., Daddi E., Dannerbauer H., 2015, *ApJ*, 798, L18
- Huško F. et al., 2026, *MNRAS*, 547, stag324
- Jiménez-Donaire M. J. et al., 2019, *ApJ*, 880, 127
- Jiménez-Donaire M. J. et al., 2023, *A&A*, 671, A3
- Jiménez E., Lagos C. d. P., Ludlow A. D., Wisnioski E., 2023, *MNRAS*, 524, 4346
- Jönsson H. et al., 2018, *AJ*, 156, 126
- Kennicutt R. C. Jr., 1983, *ApJ*, 272, 54
- Kennicutt R. C. Jr., 1989, *ApJ*, 344, 685
- Kobayashi C., Umeda H., Nomoto K., Tominaga N., Ohkubo T., 2006, *ApJ*, 653, 1145
- Kraljic K. et al., 2024, *A&A*, 682, A50
- Kroupa P., 2001, *MNRAS*, 322, 231
- Krumholz M. R., McKee C. F., 2005, *ApJ*, 630, 250
- Lacey C. G. et al., 2016, *MNRAS*, 462, 3854
- Lagos C. D. P., Lacey C. G., Baugh C. M., Bower R. G., Benson A. J., 2011, *MNRAS*, 416, 1566
- Lagos C. d. P. et al., 2015, *MNRAS*, 452, 3815
- Lagos C. d. P., Schaye J., Bahé Y., Van de Sande J., Kay S. T., Barnes D., Davis T. A., Dalla Vecchia C., 2018, *MNRAS*, 476, 4327
- Lagos C. d. P. et al., 2024, *MNRAS*, 531, 3551
- Lee J. C. et al., 2009, *ApJ*, 706, 599
- Leroy A. K., Walter F., Brinks E., Bigiel F., de Blok W. J. G., Madore B., Thornley M. D., 2008, *AJ*, 136, 2782
- Leroy A. K. et al., 2013, *AJ*, 146, 19
- Leung S.-C., Nomoto K., 2019, *PASA*, 36, e006
- Lin L. et al., 2019, *ApJ*, 884, L33
- Lin L. et al., 2022, *ApJ*, 926, 175
- Ludlow A. D., Schaye J., Schaller M., Richings J., 2019, *MNRAS*, 488, L123
- Ludlow A. D., Fall S. M., Schaye J., Obreschkow D., 2021, *MNRAS*, 508, 5114
- Ludlow A. D., Fall S. M., Wilkinson M. J., Schaye J., Obreschkow D., 2023, *MNRAS*, 525, 5614
- Martínez-Serrano F. J., Serna A., Domínguez-Tenreiro R., Mollá M., 2008, *MNRAS*, 388, 39
- McClure-Griffiths N. M., Stanimirović S., Rybarczyk D. R., 2023, *ARA&A*, 61, 19
- McGibbon R., Helly J., Schaye J., Schaller M., Vandenbroucke B., 2025, *J. Open Source Softw.*, 10, 8252
- Nagy D., Dessauges-Zavadsky M., Messa M., Richard J., Sun J., Combes F., Eyholzer Y., 2023, *A&A*, 678, A183
- Nobels F. S. J., Schaye J., Schaller M., Ploekinger S., Chaikin E., Richings A. J., 2024, *MNRAS*, 532, 3299
- Nomoto K., Tominaga N., Umeda H., Kobayashi C., Maeda K., 2006, *Nucl. Phys. A*, 777, 424
- Nomoto K., Kobayashi C., Tominaga N., 2013, *ARA&A*, 51, 457
- Orr M. E. et al., 2018, *MNRAS*, 478, 3653
- Ostriker E. C., McKee C. F., Leroy A. K., 2010, *ApJ*, 721, 975
- Park H.-J. et al., 2025, *MNRAS*, 540, 2623
- Pessa I. et al., 2021, *A&A*, 650, A134
- Pillepich A. et al., 2018, *MNRAS*, 473, 4077
- Ploekinger S., Schaye J., 2020, *MNRAS*, 497, 4857
- Ploekinger S., Richings A. J., Schaye J., Trayford J. W., Schaller M., Chaikin E., 2025, *MNRAS*, 543, 891
- Portinari L., Chiosi C., Bressan A., 1998, *A&A*, 334, 505
- Richings A. J., Schaye J., Oppenheimer B. D., 2014a, *MNRAS*, 440, 3349

Richings A. J., Schaye J., Oppenheimer B. D., 2014b, *MNRAS*, 442, 2780
 Richings A. J., Faucher-Giguère C.-A., Gurvich A. B., Schaye J., Hayward C. C., 2022, *MNRAS*, 517, 1557
 Roychowdhury S., Chengalur J. N., Begum A., Karachentsev I. D., 2009, *MNRAS*, 397, 1435
 Salpeter E. E., 1955, *ApJ*, 121, 161
 Schaller M. et al., 2024, *MNRAS*, 530, 2378
 Schaye J., 2001, *ApJ*, 562, L95
 Schaye J., 2004, *ApJ*, 609, 667
 Schaye J. et al., 2015, *MNRAS*, 446, 521
 Schaye J. et al., 2026, *MNRAS*, 548, stg375
 Schinnerer E., Leroy A. K., 2024, *ARA&A*, 62, 369
 Schmidt M., 1959, *ApJ*, 129, 243
 Schrubba A. et al., 2011, *AJ*, 142, 37
 Segovia Otero Á., Agertz O., Renaud F., Kraljic K., Romeo A. B., Semenov V. A., 2025, *MNRAS*, 538, 2646
 Semenov V. A., Kravtsov A. V., Gnedin N. Y., 2018, *ApJ*, 861, 4
 Sharon C. E., Baker A. J., Harris A. I., Thomson A. P., 2013, *ApJ*, 765, 6
 Sharon C. E. et al., 2019, *ApJ*, 879, 52
 Shen S., Wadsley J., Stinson G., 2010, *MNRAS*, 407, 1581
 Shen X. et al., 2026, *MNRAS*, 545, staf2119
 Shi Y., Helou G., Yan L., Armus L., Wu Y., Papovich C., Stierwalt S., 2011, *ApJ*, 733, 87
 Shi Y. et al., 2018, *ApJ*, 853, 149
 Somerville R. S., Popping G., Trager S. C., 2015, *MNRAS*, 453, 4337
 Sun J. et al., 2022, *AJ*, 164, 43
 Tacconi L. J., Genzel R., Sternberg A., 2020, *ARA&A*, 58, 157
 Trayford J. W. et al., 2026, *MNRAS*, 545, staf2040
 Tremmel M., Karcher M., Governato F., Volonteri M., Quinn T. R., Pontzen A., Anderson L., Bellovary J., 2017, *MNRAS*, 470, 1121
 Tremonti C. A. et al., 2004, *ApJ*, 613, 898
 Tsai J. C., Mathews W. G., 1995, *ApJ*, 448, 84
 Usero A. et al., 2015, *AJ*, 150, 115
 Walter F., Brinks E., de Blok W. J. G., Bigiel F., Kennicutt R. C., Thornley M. D., Leroy A., 2008, *AJ*, 136, 2563
 Wang J. et al., 2024, *ApJ*, 973, 15
 Wiersma R. P. C., Schaye J., Theuns T., Dalla Vecchia C., Tornatore L., 2009, *MNRAS*, 399, 574
 Wong T. et al., 2024, *ApJS*, 271, 35
 Xie L., De Lucia G., Hirschmann M., Fontanot F., Zoldan A., 2017, *MNRAS*, 469, 968
 Zanella A. et al., 2024, *A&A*, 685, A80
 Zhukovska S., Gail H. P., Tieloff M., 2008, *A&A*, 479, 453

APPENDIX A: CONVERGENCE OF THE KS RELATION

A1 Convergence with the resolution of the simulation

The top and middle panels of Fig. A1 show the KS relations for the atomic, molecular and total neutral hydrogen at $z = 0$ and 2 for the three L025 simulations listed in Table 1.

For H I, we see that the convergence at $z = 0$ is excellent, with differences between resolutions only apparent in the scatter of the relation. The resolution m5 produces the largest scatter (by 0.15 dex), consistent with Fig. 13. At $z = 2$, there are differences between the three resolutions at low surface densities of H I, $\log_{10}(\Sigma_{\text{HI}}/\text{M}_{\odot} \text{pc}^{-2}) \lesssim 1$. Below that surface density, m5 displays the steepest relationship (at $0 \lesssim \log_{10}(\Sigma_{\text{HI}}/\text{M}_{\odot} \text{pc}^{-2}) \lesssim 1$). In that regime, it follows the same slope as predicted for higher surface densities.

For H₂, we see good convergence at high surface densities. At m5 resolution, the KS slope of ≈ 1 is present at all densities $\gtrsim 0.1 \text{ M}_{\odot} \text{pc}^{-2}$. For m6 resolution, we start to see a flattening of the KS relation and a deviation from m5 resolution below

Table A1. The Spearman correlation coefficient, R , for the correlation between $\log_{10}(\tau_{\text{H}_2})$ or $\log_{10}(\tau_{\text{HI}})$ and 5 local properties listed in each row. This is shown for the L025m5, L025m6, and L025m7 simulations for galaxies with $M_{\star} > 10^9 \text{ M}_{\odot}$ and $\text{SFR} > 0$ at $z = 0$. Units used for each property are as in Fig. A3. Fig. 7 shows these values for the L200m6 run we use throughout the paper.

	$\log_{10}(\tau_{\text{HI}})$	$\log_{10}(\tau_{\text{H}_2})$
$\log_{10}(\Sigma_{\star})$	−0.66, −0.68, −0.74	0.32, 0.42, 0.35
$\log_{10}(s\Sigma_{\text{SFR}})$	−0.35, −0.34, −0.27	−0.52, −0.56, −0.44
$12 + \log_{10}(\text{O}/\text{H})$	−0.27, −0.35, −0.38	0.24, 0.32, 0.36
$\log_{10}(\Sigma_{\text{dust}})$	−0.62, −0.71, −0.77	0.26, 0.36, 0.32
$\log_{10}(\sigma_{\text{cool}})$	0.09, 0.14, 0.14	−0.08, −0.10, −0.03

$\approx 1 \text{ M}_{\odot} \text{pc}^{-2}$. For m7 resolution, this deviation occurs at higher densities, $\approx 10 \text{ M}_{\odot} \text{pc}^{-2}$. At $z = 2$, the deviations from the m5 resolution results move to lower H₂ surface densities, so that convergence is achieved already at $\gtrsim 0.1 \text{ M}_{\odot} \text{pc}^{-2}$. F. S. J. Nobels et al. (2024) used idealized galaxy discs to show that the molecular KS relation did not converge at low gas surface densities, which agrees with what we find in the fully cosmological set-up of COLIBRE.

For the total H I + H₂ KS relation, we see very similar behaviour to that obtained for the H I KS relation.

For our fiducial m6 resolution, this means that our results are converged at the typical surface densities probed by observations, $\gtrsim 10 \text{ M}_{\odot} \text{pc}^{-2}$ for H₂ and $\gtrsim 1 \text{ M}_{\odot} \text{pc}^{-2}$ for H I.

Fig. A2 shows the differences in the KS relation normalization between m6 and m5 for galaxies at $z = 0$ selected in bins of stellar mass. We find that galaxies at m6 resolution with stellar masses 10^8 – $10^9 \text{ M}_{\odot}$ can show deviations of up to 0.6 dex in the normalization of the total KS relation (right panel) relative to the m5 run. We thus decide to impose a minimum stellar mass of $10^9 \text{ M}_{\odot}$ for our analysis of the L200m6 run, but noting that there is no clear stellar mass threshold above which the KS relation is fully converged. This will depend on how star-forming/gas-rich a galaxy is, and which gas surface densities we sample. Nevertheless, we consider the stellar mass threshold above to be a reasonable compromise between probing a large dynamic range while also avoiding galaxies that are more affected by resolution.

The analysis above is for the predicted average KS relation. The convergence of the scaling relations between the molecular and atomic depletion times and the several local galaxy properties analysed in Section 3.2.1 is not guaranteed and needs to be explicitly tested. Fig. A3 shows the dependence of the H₂ and H I gas depletion times and the five properties shown in Fig. 7. In the case of τ_{HI} , Fig. A3 shows an overall good convergence between the three resolutions. This agrees with the convergence seen in the average H I KS relation in Fig. A1. The convergence is particularly good for the properties that are most strongly correlated with τ_{HI} , i.e. $\Sigma_{\star}$ and Σ_{dust} .

For τ_{H_2} we see that the m5 and m6 resolutions produce very similar correlations with the dust surface density, local sSFR, cool gas velocity dispersion and stellar surface density. The m7 resolution clearly deviates from the relations obtained at higher resolutions. The gas metallicity (top-left panel) shows that the m5 and m6 resolution exhibit significant differences, particularly in low metallicity gas ($12 + \log_{10}(\text{O}/\text{H}) \lesssim 8.7$). This may not be surprising given the different densities at which stars form in the m5 and m6 resolutions at low gas metallicity (Fig. 6). J. Schaye

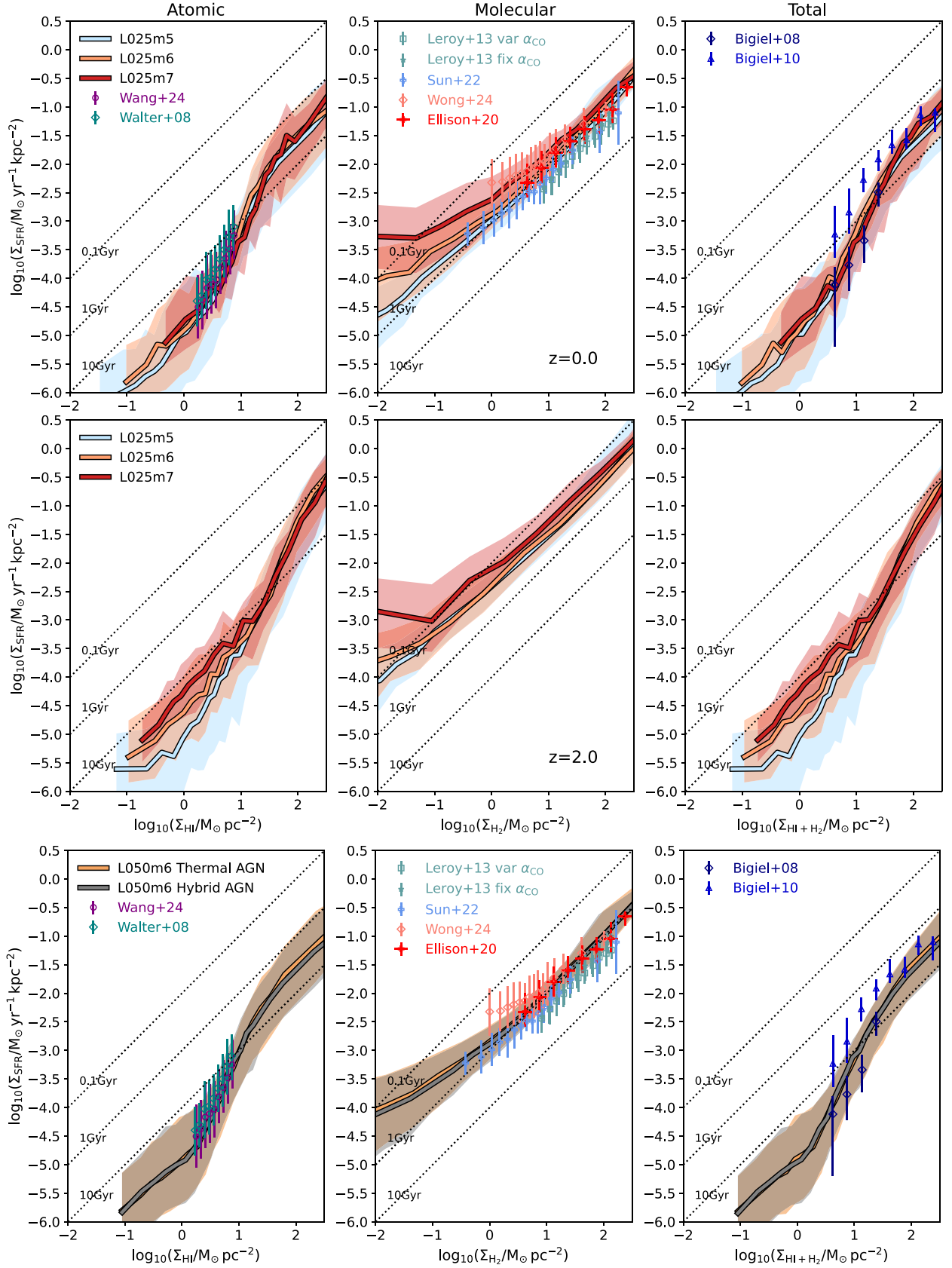


Figure A1. As Fig. 2 but for the simulations L025m5, L025m6, and L025m7, as labelled, at $z=0$ (top) and $z=2$ (middle). The bottom panel shows the $z=0$ KS relation again, but for the runs L050m6 Thermal and Hybrid, as labelled. Table 1 summarizes the runs used in this figure. Observations are only shown in the top and bottom panels.

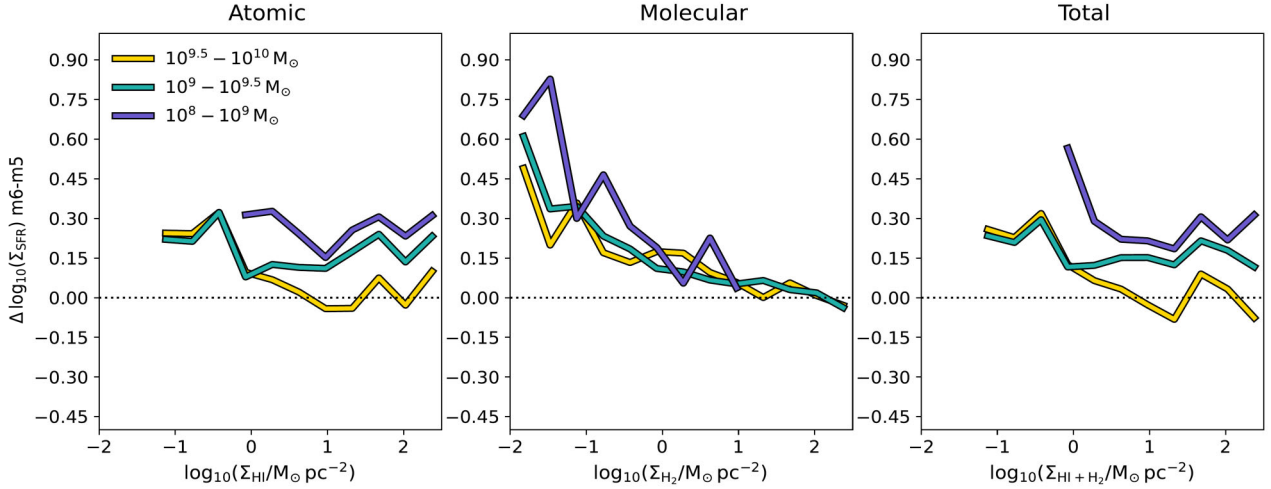


Figure A2. The difference between $\log_{10}(\Sigma_{\text{SFR}})$ in the L025m5 and L025m6 runs in bins of $\log_{10}(\Sigma_{\text{gas}})$, for galaxies at $z = 0$ selected in different ranges stellar mass, as labelled in the left panel. The differences become smaller as we move towards more massive galaxies. In the right panel, the lowest mass galaxies in the L025m6 display differences in $\log_{10}(\Sigma_{\text{SFR}})$ of up to 0.6 dex compared to L025m5.

et al. (2026) shows that the stellar mass-gas metallicity relation predicted by COLIBRE at $z = 0$ does not converge well between the runs at m5 and m6 resolution for low-mass, low-gas metallicity galaxies ($M_{\star} \lesssim 10^{9.3} M_{\odot}$).

Despite some differences found between resolutions, the strengths of the correlations shown in Fig. 7 converge quite well. We show this in Table A1. The strongest correlations in terms of Spearman’s correlation coefficient R are the same for the three resolutions. The general trend is that the m5 resolution tends to display a larger scatter in the correlations between τ_{HI} or τ_{H_2} and the five properties listed in Table A1, leading to slightly lower values of R .

A2 Convergence with the adopted AGN feedback model

The bottom panels of Fig. A1 show the KS relations predicted by the L050m6 Thermal and Hybrid AGN feedback runs (listed in Table 1). This is shown at $z = 0$ and for the same stellar mass and SFR selection. The predicted relations are very similar including the medians and scatter. The figure confirms that the adopted AGN feedback model has a negligible impact on the predicted KS relation in COLIBRE, and justifies our decision to carry out most of the analysis using runs with thermal AGN feedback only.

A3 Convergence between the methods used to measure the KS relation

Besides the method ‘Annuli on face-on galaxy’, we also implemented other methods to measure maps of the SFR and gas surface densities with the purpose of studying how well the resulting KS converges between these different methods.

Below, we summarize the methods explored:

(i) *Annuli on randomly-oriented galaxy*: This method starts by taking the simulation box XY projection, and then following the same procedure as described for the ANNULI-FACE method described in Section 2.2. We refer to this method as ANNULI-RAND.

(ii) *Regular grid on face-on galaxy*: This method also starts from the galaxies oriented face-on, which is computed as described for the ANNULI-FACE method. Then, all particles within a square of

50 kpc are selected. We then grid the space using regular square pixels of size $\Delta r \times \Delta r$. We compute for each pixel the masses, SFRs and metallicities following the ANNULI-FACE method but using only the particles that fall in the pixel. We refer to this method as GRID-FACE. Surface densities are then computed using the area of the pixel.

(iii) *Regular grid on randomly-oriented galaxy*: This method follows the description provided for GRID-FACE, but the grid is constructed over the galaxy projected on to the simulation box XY’s coordinates, as is done for ANNULI-RAND. We refer to this method as GRID-RAND.

We remind the reader that our fiducial method bins galaxies in circular apertures of $\Delta r = 1$ kpc after orienting them face-on, and with a fiducial minimum number of particles per bin of $N_{\text{min}} = 10$.

In the following, we compare different methods of measuring the KS relation at $z = 0$ in the L025m6 run to determine whether our results are robust against the choices we make about how to measure the KS relation.

The top panels of Fig. A4 shows a comparison of the 4 different methods used in this paper to measure the KS relation. For H_2 (top-middle), we find excellent agreement between methods at $\log_{10}(\Sigma_{\text{H}_2}/M_{\odot} \text{ pc}^{-2}) \gtrsim 0.5 - 1.8$, which is the regime in which observations are available. At lower and higher surface densities, we start to see some differences in the median, but they are still well within the scatter of the correlation.

For HI (top-left panel in Fig. A4), we see more systematic differences between methods, with the ones adopting a random orientation leading to a slightly higher normalization than those adopting a face-on orientation. This is naturally expected in situations where the HI disc has a larger scale-height than the disc traced by the instantaneous SFR. The latter is expected in COLIBRE given the hotter nature of HI compared to the gas particles with SFR > 0 (see Fig. 3). Regardless, the differences in normalization are within ≈ 0.2 dex, well below the intrinsic scatter of the HI KS relation in COLIBRE.

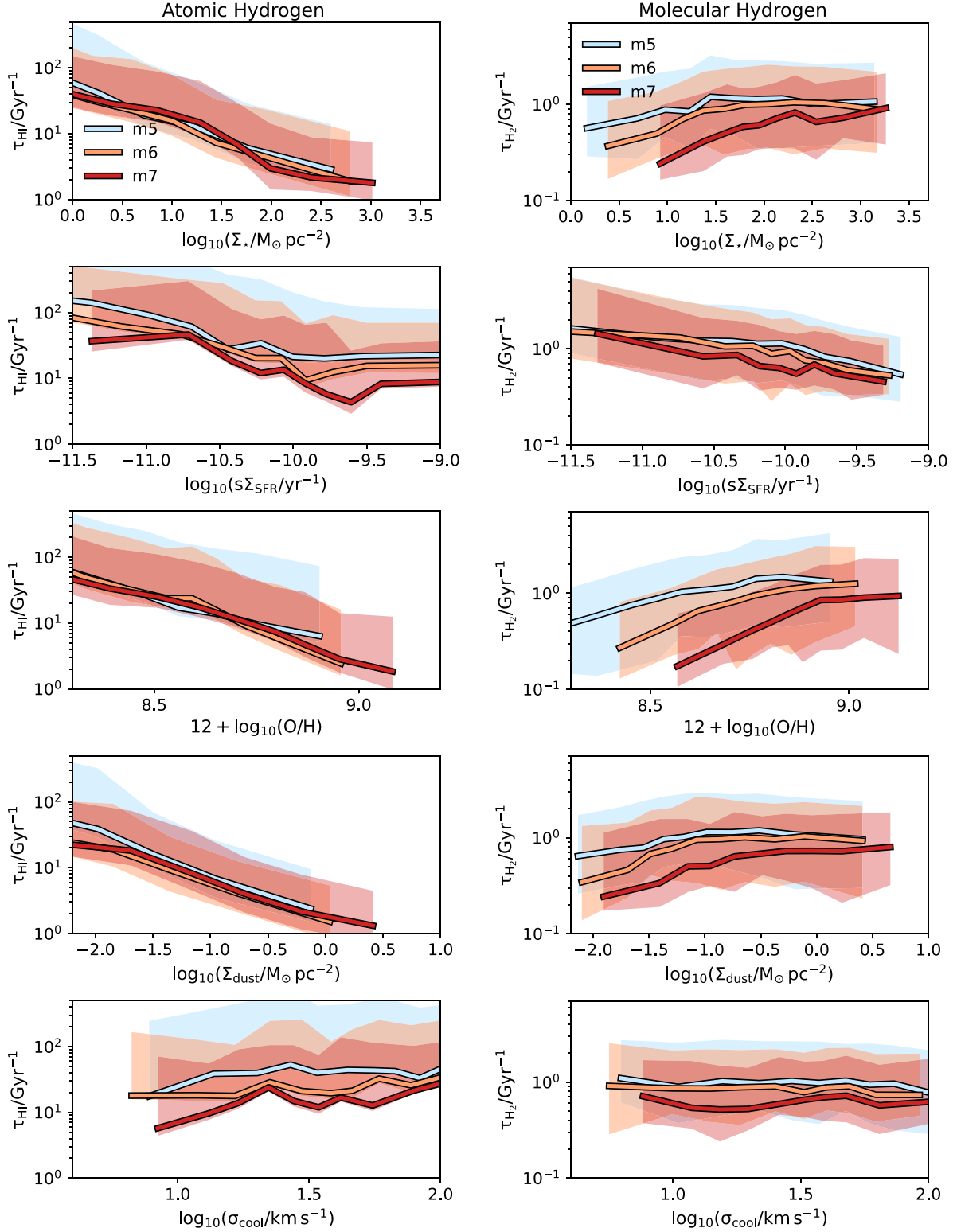


Figure A3. The relation between the depletion times of H I (left) and H₂ (right) and the 5 local properties analysed in Fig. 7—from top to bottom stellar surface density, local sSFR, gas metallicity, dust surface density, and cool gas velocity dispersion. This is shown for the L025m5, L025m6, and L025m7 as labelled in the top panels. In the three runs we include galaxies with $M_* > 10^9 M_\odot$ and $\text{SFR} > 0$ at $z = 0$.

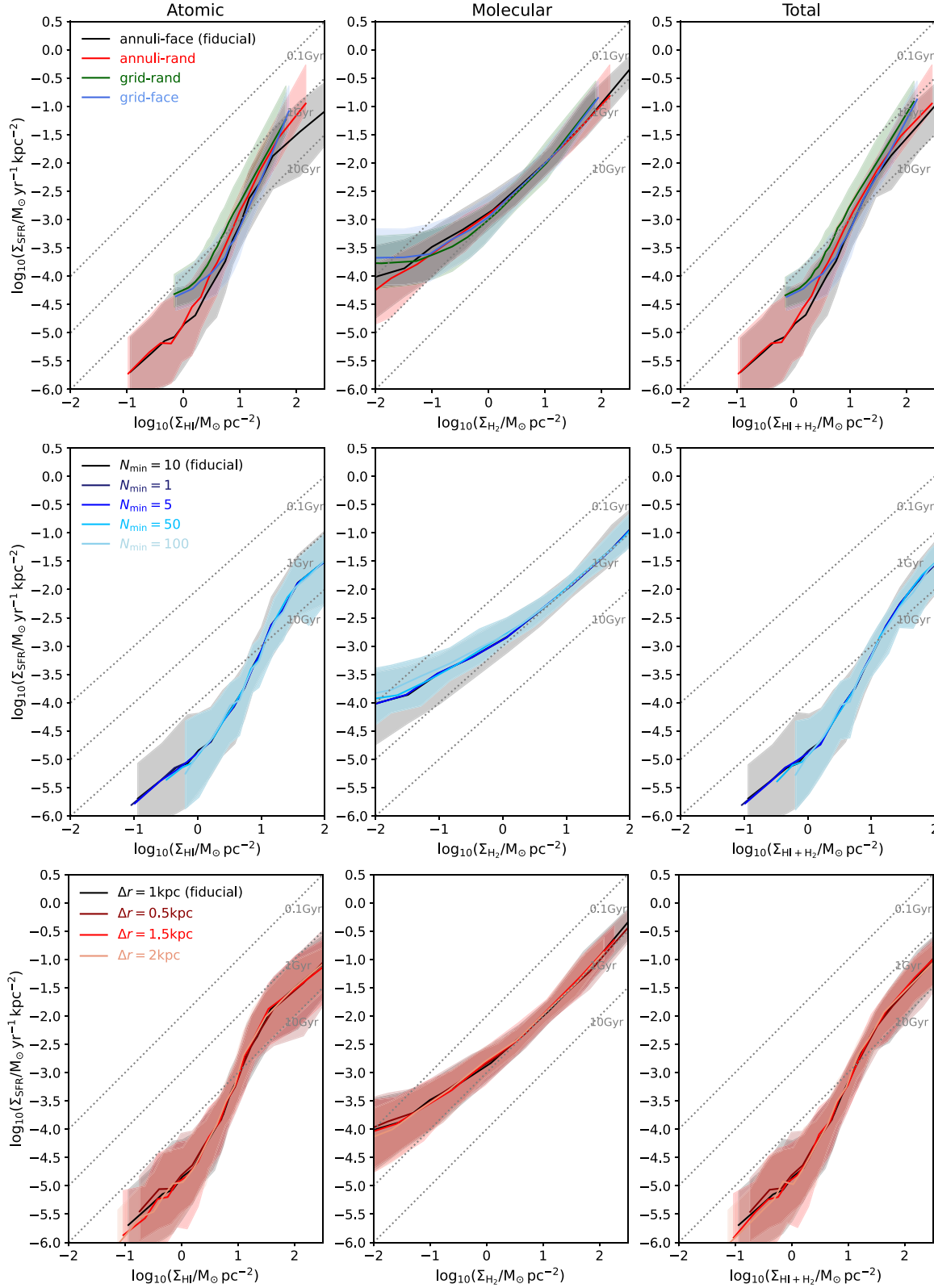


Figure A4. The H I, H₂, and H I + H₂ KS relations, as labelled at the top, at $z = 0$ in the L025m6 run of Table 1, measured using different methods (top); using our fiducial ANNULI-FACE method but adopting different thresholds for the minimum number of gas particles in an annulus (middle); and adopting different widths for the annuli (bottom). In all panels lines show medians, while the shaded regions indicate the 16th–84th percentile ranges. Not not overwhelm the figure, in the middle panel we only show the 16th–84th percentile ranges for $N_{\text{min}} = 10$ and $= 100$, while for the other panels we show medians and percentile ranges for all the samples. We find excellent convergence in the predicted KS relation against methods and the specific choices of each method.

For the total $\text{H I} + \text{H}_2$ (top-right panel of Fig. A4), the systematic differences between methods are smaller than those seen for H I , but go in the same direction, with the methods using face-on orientation producing a slightly lower normalization of the KS relation than those using random orientations.

The middle and lower panels of Fig. A4 show the KS relations obtained for COLIBRE using the ANNULI-FACE method, but varying Δr (the width of the circular bins) and $N_{\min}$ (the minimum number of particles required per circular bin), as labelled. In both cases we see excellent convergence in the shape and scatter of the KS relation in all neutral hydrogen phases. The main difference we see is that increasing $N_{\min}$ leads to a poorer sampling of the low surface density regime, but without introducing a significant bias.

Overall, Fig. A4 demonstrate that the predicted KS relation in COLIBRE at the m6 resolution is robust against variants in the methods used to measure it.

APPENDIX B: ADDITIONAL VISUALIZATION EXAMPLES

Fig. B1 shows examples of star-forming galaxies (top) and passive galaxies (bottom) at $z = 0$ in the L025m6 run. The stellar mass and its hierarchy (whether it is a central or a satellite galaxy) are shown in the top panels. In general and at this resolution, we see that star-forming galaxies sample H_2 surface densities in excess of $\log_{10}(\Sigma_{\text{H}_2}/\text{M}_{\odot} \text{pc}^{-2}) \gtrsim 1.5$, while galaxies below the main sequence rarely have H_2 gas that reaches such high densities.

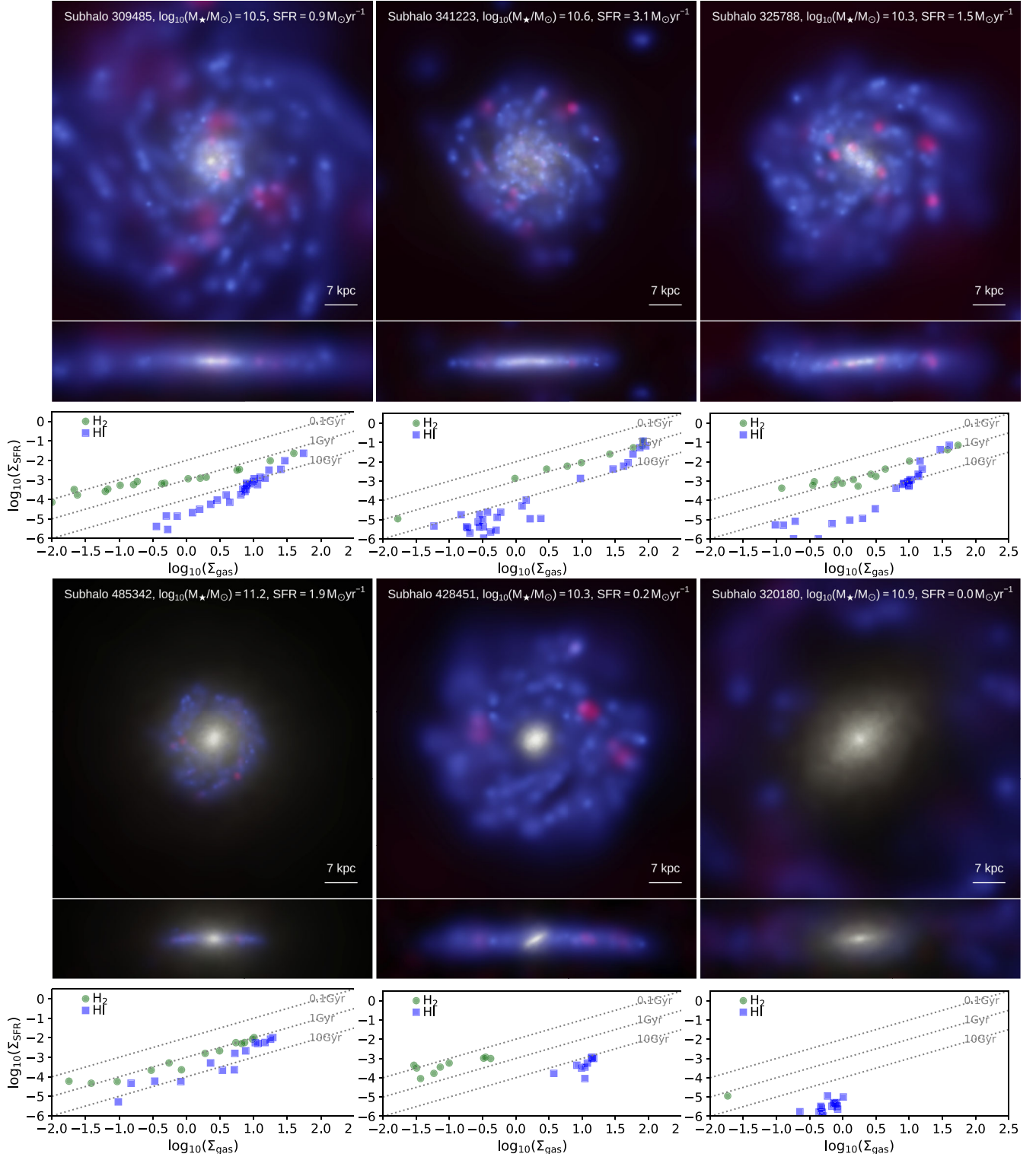


Figure B1. As Fig. 1 but for another 6 galaxies at $z = 0$ in the L025m6 run. The top row shows examples of star-forming galaxies, while the bottom row shows examples of galaxies below the main sequence but that still have some gas left to form stars.

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